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THREE MILLISECOND PULSARS IN *FERMI* LAT UNASSOCIATED BRIGHT SOURCES

S. M. RANSOM¹, P. S. RAY², F. CAMILO³, M. S. E. ROBERTS⁴, Ö. ÇELİK^{5,6,7}, M. T. WOLFF², C. C. CHEUNG^{8,21}, M. KERR⁹,
T. PENNUCCI¹⁰, M. E. DECESAR^{5,11}, I. COGNARD¹², A. G. LYNE¹³, B. W. STAPPERS¹³, P. C. C. FREIRE¹⁴, J. E. GROVE²,
A. A. ABDO⁸, G. DESVIGNES^{15,16}, D. DONATO^{6,11}, E. C. FERRARA⁵, N. GEHRELS⁵, L. GUILLEMOT¹⁴, C. GWON², A. K. HARDING⁵,
S. JOHNSTON¹⁷, M. KEITH¹⁷, M. KRAMER^{13,14}, P. F. MICHELSON⁹, D. PARENT^{18,21}, P. M. SAZ PARKINSON¹⁹, R. W. ROMANI⁹,
D. A. SMITH²⁰, G. THEUREAU¹², D. J. THOMPSON⁵, P. WELTEVREDE¹³, K. S. WOOD², AND M. ZIEGLER¹⁹

¹ National Radio Astronomy Observatory (NRAO), Charlottesville, VA 22903, USA; sransom@nrao.edu

² Space Science Division, Naval Research Laboratory, Washington, DC 20375, USA; Paul.Ray@nrl.navy.mil

³ Columbia Astrophysics Laboratory, Columbia University, New York, NY 10027, USA

⁴ Eureka Scientific, Oakland, CA 94602, USA

⁵ NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA

⁶ Center for Research and Exploration in Space Science and Technology (CREST) and NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA

⁷ Department of Physics and Center for Space Sciences and Technology, University of Maryland Baltimore County, Baltimore, MD 21250, USA

⁸ National Research Council Research Associate, National Academy of Sciences, Washington, DC 20001, USA

⁹ W. W. Hansen Experimental Physics Laboratory, Kavli Institute for Particle Astrophysics and Cosmology, Department of Physics and SLAC National Accelerator Laboratory, Stanford University, Stanford, CA 94305, USA

¹⁰ Department of Astronomy, University of Virginia, Charlottesville, VA 22904, USA

¹¹ Department of Physics and Department of Astronomy, University of Maryland, College Park, MD 20742, USA

¹² Laboratoire de Physique et Chimie de l'Environnement, LPCE UMR 6115 CNRS, F-45071 Orléans Cedex 02, and Station de radioastronomie de Nançay, Observatoire de Paris, CNRS/INSU, F-18330 Nançay, France

¹³ Jodrell Bank Centre for Astrophysics, School of Physics and Astronomy, The University of Manchester, M13 9PL, UK

¹⁴ Max-Planck-Institut für Radioastronomie, Auf dem Hügel 69, 53121 Bonn, Germany

¹⁵ Department of Astronomy, University of California, Berkeley, CA 94720-3411, USA

¹⁶ Radio Astronomy Laboratory, University of California, Berkeley, CA 94720, USA

¹⁷ CSIRO Astronomy and Space Science, Australia Telescope National Facility, Epping, NSW 1710, Australia

¹⁸ College of Science, George Mason University, Fairfax, VA 22030, USA

¹⁹ Santa Cruz Institute for Particle Physics, Department of Physics and Department of Astronomy and Astrophysics, University of California at Santa Cruz, Santa Cruz, CA 95064, USA

²⁰ Université Bordeaux 1, CNRS/IN2p3, Centre d'Études Nucléaires de Bordeaux Gradignan, 33175 Gradignan, France

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ABSTRACT

We searched for radio pulsars in 25 of the non-variable, unassociated sources in the *Fermi* LAT Bright Source List with the Green Bank Telescope at 820 MHz. We report the discovery of three radio and γ -ray millisecond pulsars (MSPs) from a high Galactic latitude subset of these sources. All of the pulsars are in binary systems, which would have made them virtually impossible to detect in blind γ -ray pulsation searches. They seem to be relatively normal, nearby (≤ 2 kpc) MSPs. These observations, in combination with the *Fermi* detection of γ -rays from other known radio MSPs, imply that most, if not all, radio MSPs are efficient γ -ray producers. The γ -ray spectra of the pulsars are power law in nature with exponential cutoffs at a few GeV, as has been found with most other pulsars. The MSPs have all been detected as X-ray point sources. Their soft X-ray luminosities of $\sim 10^{30}$ – 10^{31} erg s⁻¹ are typical of the rare radio MSPs seen in X-rays.

Key words: pulsars: general – pulsars: individual (J0614–3329, J1231–1411, J2214+3000)

1. INTRODUCTION

Before the launch of the *Fermi* Gamma-ray Space Telescope, the only pulsars with definitive detections in γ -rays (using EGRET on the *Compton Gamma Ray Observatory*) were young and very energetic ($\dot{E} > 10^{36}$ erg s⁻¹) or nearby older systems ($\dot{E} > 10^{34}$ erg s⁻¹; Thompson 2004). A possible detection of pulsed γ -rays from the energetic millisecond pulsar (MSP) J0218+4232 (Kuiper et al. 2000) sparked interest in modeling MSP γ -ray emission (e.g., Zhang & Cheng 2003; Harding et al. 2005) and encouraged one group (Story et al. 2007) to predict that many new MSPs might be detected in γ -rays or discovered in radio follow-up of unidentified *Fermi* sources.

The launch of *Fermi* and the extraordinary sensitivity of the Large Area Telescope (LAT; Atwood et al. 2009) confirmed those predictions of γ -ray-bright MSPs with detections of eight relatively normal radio MSPs using only the first few months of

Fermi events (Abdo et al. 2009a). Those MSPs were detected via the folding of γ -rays modulo the known spin and orbital ephemerides from radio timing campaigns (Smith et al. 2008).

In order to best utilize radio telescope time to search either for radio counterparts to new γ -ray-selected pulsars or to search blindly for radio pulsations from γ -ray sources that might contain pulsars, we formed the Pulsar Search Consortium (PSC), a group of approximately 20 LAT team members and/or pulsar experts associated with large radio telescopes around the world. This Letter describes one of the PSC's first programs, which used the Green Bank Telescope (GBT) to search 25 unassociated sources from the *Fermi* LAT Bright Source List (BSL; Abdo et al. 2009b).

2. OBSERVATIONS AND DATA ANALYSIS

We selected 25 sources from the *Fermi* LAT Bright Source List that were (1) unassociated with known pulsars or active galactic nuclei (AGNs), (2) unassociated with X-ray counterparts that had been previously deeply searched for radio

²¹ Current address: Naval Research Laboratory, Washington, DC 20375, USA.

pulsations (e.g., IC443; F. Camilo et al. 2011, in preparation), (3) statistically non-variable, and (4) at declinations $> -35^\circ$. We observed each of the sources²² for approximately 45–50 minutes using the prime focus receiver at the GBT centered at 820 MHz with 200 MHz of bandwidth. The GBT pointings, all taken between 2009 July and October, were actually centered on the positions from an internal LAT source list using nine months of sky-survey data prepared in a similar fashion to the *Fermi* LAT First Source Catalog (i.e., “1FGL”; Abdo et al. 2010). The individual GBT pointings at 820 MHz had FWHM = $0^\circ.25$ and covered either all or a substantial fraction of the 95% error regions for the vast majority of the sources.

We sampled the summed power from two polarizations in 2048 frequency channels with 8-bits every $61.44 \mu\text{s}$ using the GUPPI pulsar backend.²³ Each pointing generated approximately 100 GB of data, which were recorded to hard drives for processing off-site. The 820 MHz center frequency was chosen as a compromise between the competing effects of sky temperature (from the Galactic synchrotron background) and beam size as well as steep pulsar spectra and the effects of interstellar dispersion and scattering.

For a pulsar with a pulse width of $\sim 10\%$ of the period, the search sensitivity was approximately $0.06 \times (29 \text{ K} + T_{\text{sky}})/(32 \text{ K}) \text{ mJy}$, where T_{sky} is the contribution at 820 MHz of the Galactic synchrotron background. The majority of the sources (17 of them) were within three degrees of the Galactic plane where $T_{\text{sky}} \sim 12\text{--}40 \text{ K}$, although two were very near the Galactic center with $T_{\text{sky}} \sim 100\text{--}150 \text{ K}$. Eight of the sources were well off the Galactic plane ($|b| > 5^\circ$) and had $T_{\text{sky}} \sim 3\text{--}10 \text{ K}$. For those sources, our search sensitivity was $0.06\text{--}0.08 \text{ mJy}$ for normal pulsars at all reasonable dispersion measures (DMs) and MSPs up to $\text{DM} \sim 100 \text{ pc cm}^{-3}$. In general, the observations were factors of 2–12 deeper than the best pulsar surveys that have previously covered these regions (e.g., Manchester et al. 2001; Cordes et al. 2006).

We processed the data, after de-dispersing into ~ 9000 DMs over the range $0\text{--}1055 \text{ pc cm}^{-3}$, using both acceleration searches (to improve sensitivity to pulsars in binary systems) and single pulse searches (to provide sensitivity to pulsars with sporadic or giant-pulse-like emission) using standard tools found in PRESTO²⁴ (Ransom et al. 2002). No new pulsar-like signals were found in any of the low Galactic latitude ($|b| < 5^\circ$) sources. However, four new MSPs were detected amid the eight high Galactic latitude sources, in 0FGLs J0614.3–3330, J1231.5–1410, J2214.8+3002, and J2302.9+4443. The MSP in 0FGL J2302.9+4443 was detected first in an independent PSC survey by the Nançay telescope and will be reported elsewhere (Cognard et al. 2011). The rest of this Letter details the properties of the other three MSPs.

2.1. The New MSPs

The first two pulsars detected, J2214+3000 and J1231–1411, were undergoing substantial accelerations due to orbital motion during the discovery observations. PSR J0614–3329 was initially uncovered in an unaccelerated search, although orbital

motion was detected in the discovery observations via a more precise timing analysis. The fact that all three MSPs were in unknown binaries demanded a radio timing program to determine precise orbital parameters and constrain their astrometric positions before detailed γ -ray timing and analysis could commence.

Each MSP was observed with several different observing setups at the GBT, the Lovell Telescope at Jodrell Bank, and the Nançay radio telescope. The Arecibo telescope also observed PSR J2214+3000 several times. At the GBT, GUPPI was used with bandwidths of 100, 200, and 800 MHz centered at 350, 820, and 1500 MHz, respectively. At Jodrell Bank and Nançay, observations were made with bandwidths of 200–300 MHz centered near 1400 MHz. Standard radio timing procedures were used (Lorimer & Kramer 2005) and the orbital parameters were fit to high precision with TEMP02 (Hobbs et al. 2006).

After several months of radio timing and using trial X-ray positions based on point sources from *Swift*, *XMM-Newton*, and/or *Chandra*, γ -ray pulsations from each MSP were detected with the LAT using orbitally demodulated events (see also Section 2.2). We determined average γ -ray pulse times of arrival (TOAs) using the maximum likelihood γ -ray timing techniques described by Ray et al. (2010) after integrating source photons modulo the predicted pulse period between 22 and 36 days per TOA. The resulting joint timing solutions, using both radio and γ -ray TOAs, as well as the derived physical parameters of the MSPs, are presented in Table 1. The radio and γ -ray pulse profiles are shown in Figure 1.

The three new pulsars, besides being three of the brightest γ -ray MSPs in the sky, appear to be relatively normal, nearby ($\leq 2 \text{ kpc}$), radio MSPs, with $\sim 3 \text{ ms}$ spin periods, surface magnetic field strengths of $(2\text{--}3) \times 10^8 \text{ G}$, and spin-down luminosities of $\sim 2 \times 10^{34} \text{ erg s}^{-1}$. J0614–3329 and J1231–1411 have orbital periods of 53.6 and 1.9 days, respectively, with companions of mass $\sim 0.2\text{--}0.3 M_\odot$, consistent with the orbital period–white dwarf mass relation of Rappaport et al. (1995). J2214+3000 is a so-called black-widow system with a very low mass companion ($\sim 0.02 M_\odot$) and likely timing irregularities, similar to pulsars B1957+20 (Fruchter et al. 1988), J2051–0827 (Stappers et al. 1996), and J0610–2100 (Burgay et al. 2006), and only the fourth such system known in the Galactic disk. While we currently have no evidence for radio eclipses from the pulsar (at least at frequencies $\geq 1.4 \text{ GHz}$), its formation was likely similar to that of those other systems (e.g., King et al. 2005).

2.2. γ -ray Analysis

The *Fermi* LAT is sensitive to γ -rays with energies $0.02\text{--}300 \text{ GeV}$ (Atwood et al. 2009). The *Fermi* LAT sky-survey data set used here for spectral analysis spans from 2008 August 4 to 2010 February 4. We selected “Pass 6 Diffuse” class events—i.e., events passing the most stringent background rejection cuts—with energies above 0.1 GeV and rejected events with zenith angles $> 105^\circ$ to limit contamination from γ -rays from the Earth’s limb. We used “Pass6 v3” instrument response functions (IRFs).

The γ -ray light curves shown in Figure 1 are constructed from events as described above, although using seven additional months of data through 2010 September 14, and with energy and radius cuts to optimize the signal to noise for each pulsar. The energy and radius cuts used for PSRs J0614–3329, J1231–1411, and J2214+3000 are $(1^\circ, 0.35 \text{ GeV})$, $(1.2^\circ, 0.35 \text{ GeV})$, and $(1^\circ, 0.7 \text{ GeV})$, respectively.

²² The LAT Bright Sources observed with the GBT were 0FGLs J0614.3–3330, J1231.5–1410, J1311.9–3419, J1653.4–0200, J1741.4–3046, J1746.0–2900, J1801.6–2327, J1805.3–2138, J1814.3–1739, J1821.4–1444, J1834.4–0841, J1836.1–0727, J1839.0–0549, J1844.1–0335, J1848.6–0138, J1855.9+0126, J1900.0+0356, J1911.0+0905, J1923.0+1411, J2001.0+4352, J2027.5+3334, J2110.8+4608, J2214.8+3002, J2302.9+4443, and J2339.8–0530.

²³ <https://safe.nrao.edu/wiki/bin/view/CICADA/NGNPP>

²⁴ <http://www.cv.nrao.edu/~sransom/presto/>

Table 1
Parameters for the New MSPs

Parameter	PSR J0614–3329	PSR J1231–1411	PSR J2214+3000
<i>Fermi</i> BSL association (0FGL)	J0614.3–3330	J1231.5–1410	J2214.8+3002
<i>Fermi</i> 1 year source (1FGL)	J0614.1–3328	J1231.1–1410	J2214.8+3002
Timing parameters			
Right Ascension (R.A., J2000)	06 ^h 14 ^m 10 ^s .3478(3)	12 ^h 31 ^m 11 ^s .3132(7)	22 ^h 14 ^m 38 ^s .8460(1)
Declination (decl., J2000)	–33°29′54″.118(4)	–14°11′43″.63(2)	+30°00′38″.234(4)
Proper motion in R.A. (mas yr ^{–1})	...	–1.0(2) × 10 ²	...
Proper motion in decl. (mas yr ^{–1})	...	–3(4) × 10 ¹	...
Pulsar period (ms)	3.148669579439(9)	3.683878711077(3)	3.119226579079(4)
Pulsar frequency (Hz)	317.5944552995(9)	271.4530196103(2)	320.5922925597(4)
Frequency derivative (Hz s ^{–1})	–1.77(7) × 10 ^{–15}	–1.68(1) × 10 ^{–15}	–1.44(3) × 10 ^{–15}
Frequency 2nd deriv. (Hz s ^{–2})	1.7(4) × 10 ^{–23}
Reference epoch (MJD)	55100	55100	55100
Dispersion measure (pc cm ^{–3})	37.049(1)	8.090(1)	22.557(1)
Orbital period (days)	53.5846127(8)	1.860143882(9)	0.416632943(5)
Projected semimajor axis (lt-s)	27.638787(2)	2.042633(3)	0.0590800(9)
Orbital eccentricity	0.0001801(1)	4(3) × 10 ^{–6}	< 2 × 10 ^{–4}
Longitude of periastron (deg)	15.92(4)	3.2(4) × 10 ²	...
Epoch of periastron (MJD)	55146.821(7)	55016.8(2)	...
Epoch of ascending node (MJD)	55094.137854(2)
Span of timing data (MJD)	54683–55422	54683–55430	54683–55415
Number of γ -ray TOAs	24	32	20
rms γ -ray TOA residual (μ s)	99.1	24.4	110.2
Number of radio TOAs	328	136	437
rms radio TOA residual (μ s)	7.1	9.3	5.0
Derived parameters			
Mass function (M_{\odot})	0.007895133(3)	0.00264460(2)	1.2755(1) × 10 ^{–6}
Min companion Mass (M_{\odot})	≥ 0.28	≥ 0.19	≥ 0.014
Galactic longitude (deg)	240.50	295.53	86.86
Galactic latitude (deg)	–21.83	48.39	–21.67
DM-derived distance (kpc)	1.9	0.4	1.5
Flux density at 820 MHz (mJy)	1.5	0.4	2.1
Surface magnetic field (10 ⁸ G)	2.4	2.9	2.1
Characteristic age (Gyr)	2.8	2.6	3.5
Spin-down lumin. \dot{E} (10 ³⁴ erg s ^{–1})	2.2	1.8	1.8
γ -ray spectral fit parameters			
K (10 ^{–11} ph cm ^{–2} s ^{–1} MeV ^{–1})	2.12 ± 0.10 ± 0.13	2.62 ± 0.16 ± 0.18	0.94 ± 0.11 ± 0.05
Spectral index Γ	1.44 ± 0.05 ± 0.07	1.40 ± 0.07 ± 0.05	1.44 ± 0.13 ± 0.11
E_{cutoff} (GeV)	4.49 ± 0.54 ^{+1.38} _{–0.84}	2.98 ± 0.33 ^{+0.43} _{–0.29}	2.53 ± 0.50 ^{+0.43} _{–0.29}
F_{100} (10 ^{–8} ph cm ^{–2} s ^{–1})	9.52 ± 0.46 ± 0.45	10.57 ± 0.62 ± 0.39	3.83 ± 0.44 ± 0.06
G_{100} (10 ^{–11} erg cm ^{–2} s ^{–1})	10.86 ± 0.35 ± 1.06	10.33 ± 0.35 ± 0.87	3.32 ± 0.21 ± 0.24
TS	5270.3	4798.4	958.0
$\Delta\text{TS}_{\text{cutoff}}$	184.7	203.3	63.3
η (%)	210	11	49
X-ray parameters			
X-ray source R.A. (J2000)	06 ^h 14 ^m 10 ^s .3(3)	12 ^h 31 ^m 11 ^s .3(4)	22 ^h 14 ^m 38 ^s .84(3)
X-ray source decl. (J2000)	–33°29′54″(5)	–14°11′43″(6)	30°00′38″.2(6)
BB temperature (keV)	0.23(5)	0.21(5)	0.25(4)
$F_{\text{BB},0.5-8\text{keV}}$ (10 ^{–14} erg cm ^{–2} s ^{–1})	8.7 ^{+3.4} _{–3.9}	15 ^{+5.3} _{–7.4}	2.9 ^{+0.6} _{–0.7}
$L_{\text{BB},0.5-8\text{keV}}$ (10 ³⁰ erg s ^{–1})	38 ⁺¹⁵ _{–17}	2.9 ^{+1.0} _{–1.4}	7.8 ^{+1.6} _{–1.9}

Notes. Numbers in parentheses represent 2σ uncertainties in the last digit as determined by TEMPO2 using the DE405 Solar System Ephemeris for the timing parameters and 1σ uncertainties for the other parameters. The time system used is Barycentric Dynamical Time (TDB). Minimum companion masses were calculated assuming a pulsar mass of $1.4 M_{\odot}$. The DM distances were estimated using the NE2001 Galactic electron density model and likely have $\sim 20\%$ uncertainties (Cordes & Lazio 2002). The gamma-ray spectral parameters are from fits of exponentially cutoff power laws as described in Section 2.2. F_{100} and G_{100} give the integrated photon or energy flux above 0.1 GeV, respectively, while the last two parameters are gamma-ray detection significance of the source and significance of an exponential cutoff (as compared to a simple power law), where the approximate Gaussian significance is given by $\sim \sqrt{\text{TS}}$, and TS is the Test Statistic $\text{TS} = 2\Delta \log(\text{likelihood})$ between models with and without the source. The first errors are statistical and the second errors are systematic errors calculated from the bracketing IRFs. The γ -ray efficiency, $\eta = L_{\gamma} f_{\Omega} / \dot{E} = 4\pi D^2 G_{100} / \dot{E}$, assumes a beaming correction factor, $f_{\Omega} = 1$. The X-ray results are from *Swift* for PSRs J0614–3329 and J1231–1411, and *Chandra* for PSR J2214+3000. The quoted positional uncertainty for J2214+3000 is dominated by a systematic error of 0′.6 in the absolute *Chandra* astrometry (statistical error = 0′.2). X-ray fluxes and luminosities from 0.5 to 8 keV, $F_{\text{BB},0.5-8\text{keV}}$, and $L_{\text{BB},0.5-8\text{keV}}$, are based on blackbody fits to the point source counts using $kT = 0.16$ keV.

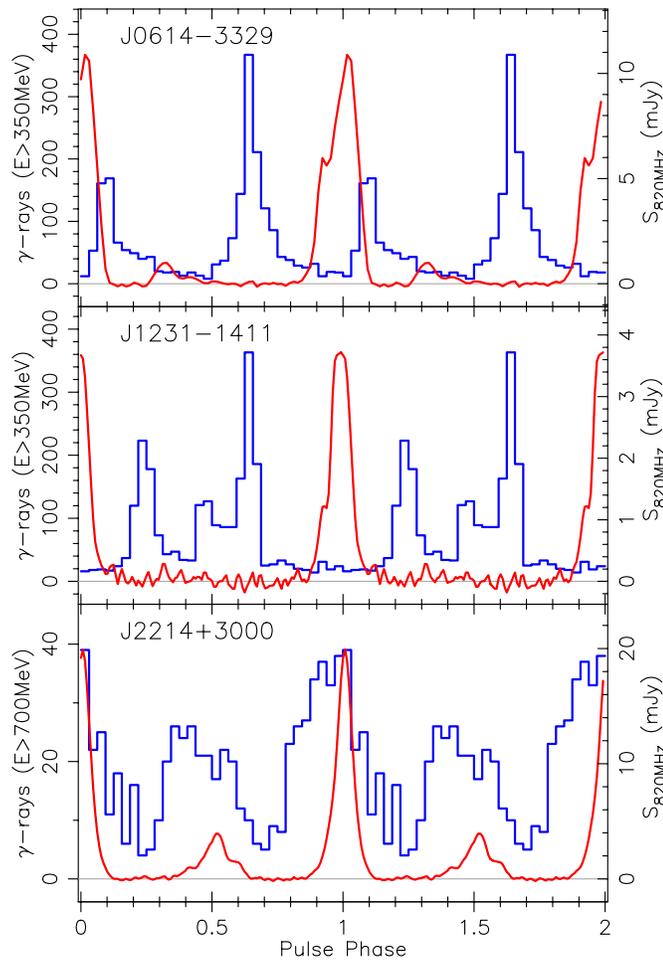


Figure 1. Radio and γ -ray pulse profiles for the three new MSPs. The red lines correspond to the 820 MHz discovery pulse profiles from the GBT with the 820 MHz flux density scale on the right. The blue lines are the *Fermi* pulse profiles with the photon counts and low energy cut used for the γ -ray selections listed on the left. There are 2043, 2341, and 621 photons in the γ -ray profiles for PSRs J0614–3329, J1231–1411, and J2214+3000, respectively.

We derived the γ -ray spectrum of each pulsar using a maximum-likelihood method implemented in the *LAT Science Tool* *gtlike*.²⁵ We analyzed a region of 10° radius centered on the radio position of each pulsar and modeled each region by including all sources from the 1FGL (Abdo et al. 2010) within 17° of the pulsar along with Galactic and isotropic diffuse emission (models *gll_iem_v02* and *isotropic_iem_v02*, respectively²⁶). The power-law spectral parameters for all sources within 10° of the pulsar and a normalizing scale factor for the diffuse emission spectrum were allowed to be free in the fit.

We modeled the spectrum of each MSP using a power law with an exponential cutoff where the three parameters, the differential flux K , the photon index Γ , and the cutoff energy E_{cutoff} , were allowed to vary in the fit. The phase-averaged γ -ray spectra obtained for each pulsar are shown in Figure 2 and the spectral parameters are given in Table 1. The uncertainty in the *LAT* effective area is estimated to be $\leq 5\%$ near 1 GeV, 10% below 0.1 GeV, and 20% over 10 GeV. The resulting systematic errors on the three spectral parameters, propagated from the uncertainties on the *LAT* effective area, were calculated using a set of “modified IRFs” bracketing the nominal (Pass6 v3) one.

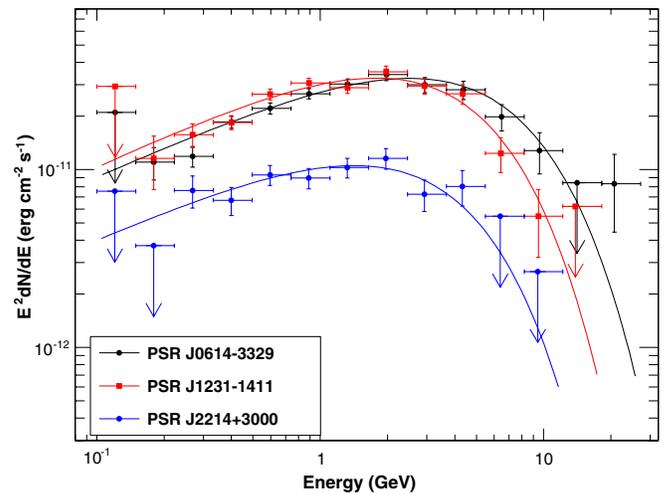


Figure 2. γ -ray spectra for the three MSPs. The flux points on the curve were obtained from independent fits in each energy bin, as explained in the text. The curves represent the spectrum fit with a simple exponentially cutoff power law in the full energy range of 0.1–100 GeV.

We verified the significance of the exponential cutoff in each spectrum with a likelihood ratio test (Mattox et al. 1996). A simple power-law model is rejected significantly for all three pulsars relative to an exponentially cutoff power law, as indicated by the test statistic $\Delta\text{TS}_{\text{cutoff}}$ listed in Table 1 for the addition of one free parameter.

The flux points in Figure 2 were obtained by repeating the likelihood analysis in each energy band, assuming a power-law spectrum with a photon index fixed at 2 and a free flux normalization parameter for all sources.

2.3. X-ray Analysis

To search for X-ray counterparts, we observed the field of each pulsar with the *Swift*-XRT (Burrows et al. 2005) with exposures ranging from 2.6 to 15.9 ks. For J0614–3329 and J1231–1411, we detected X-ray point sources at the locations of the pulsars. In the *Swift* observation of PSR J2214+3000, there is no significant source at the location of the pulsar, however we also obtained a *Chandra* ACIS-I observation of this region and detected the pulsar with it. A detailed analysis of this *Chandra* observation will be presented elsewhere.

The X-ray counterparts to the MSPs are soft sources and were fit to blackbody spectra using XSPEC, fixing the absorption values to zero. The resultant temperatures for the three X-ray sources were in the range ~ 0.21 – 0.25 keV, which in the case of J1231–1411, is consistent with that derived from deeper *Suzaku* data (Maeda et al. 2011). Allowing for additional Galactic absorption fixed to the values of $(3.4$ – $5.7) \times 10^{20} \text{ cm}^{-2}$ from Kalberla et al. (2005), we found negligible differences in the fitted parameters. The *Swift* and *Chandra* positional localizations, as well as approximate fluxes and luminosities from the blackbody fits, are in Table 1.

For PSR J1231–1411 we were able to do more detailed spectral analysis using *XMM-Newton*. On 2009 July 15 *XMM-Newton* observed the field of this as-yet unidentified BSL source with all three EPIC instruments. Data from each instrument were analyzed utilizing the Science Analysis System software version 10.0.0 and the calibration update of 2010 May 8. We filtered the data for bad events and excluded times of high particle background, yielding 29.4 ks and 29.5 ks of good time for the EPIC-MOS1 and -MOS2 instruments and 24.1 ks

²⁵ <http://fermi.gsfc.nasa.gov/ssc/data/analysis/documentation/>

²⁶ <http://fermi.gsfc.nasa.gov/ssc/data/access/lat/BackgroundModels.html>

for EPIC-PN. All three instruments utilized Full Frame mode so none had sufficient time resolution to allow searches for X-ray pulsations.

The *XMM-Newton* X-ray images of the J1231–1411 field reveal a relatively isolated, moderately bright, point source that we name XMMUJ123112–141146 at the best radio position of the pulsar to the accuracy of the EPIC instruments. We generated spectra of this source by extracting events from around the pulsar in 35'' regions for the MOS images and a 25'' region (due to a chip gap) for the PN image. This yielded spectra consisting of 853, 812, and 2164 events from the MOS1, MOS2, and PN cameras, respectively. We group the counts with at least 30 events per spectral bin for the MOS spectra and 45 events per bin for the PN spectra using the FTOOL `grppha`. Background spectra were extracted from nearby regions 100'' and 55'' in radius from the same CCD chips for the MOS and PN instruments, respectively.

Using XSPEC (v12.6.0) we simultaneously fit the MOS1, MOS2, and PN spectra in the energy range 0.4–3.0 keV. An absorbed power-law model is formally acceptable with a reduced $\chi^2 = 0.977$ (61 dof) but with an extremely steep photon index of $\Gamma = 4.23^{+0.41}_{-0.38}$ (90% confidence) and column density of $n_{\text{H}} = (1.8^{+0.6}_{-0.5}) \times 10^{21} \text{ cm}^{-2}$ (90% confidence). Such a column density is significantly higher than that expected for this direction based on galactic surveys ($n_{\text{H}} = 3.45 \times 10^{20} \text{ cm}^{-2}$; Kalberla et al. 2005). A fit to these data utilizing a model of an absorbed neutron star non-magnetic hydrogen atmosphere (*phabs* \times *nsatmos*; Heinke et al. 2006) with the neutron star mass and radius held fixed at 1.4 M_{\odot} and 10 km, and the source distance fixed at the DM value of 0.4 kpc yields a reduced $\chi^2 = 1.41$ (61 dof). In this model there are significant residuals above 1.5 keV so we add a power-law component with photon index fixed at 1.8 and obtain an improved fit with $\chi^2 = 1.09$ (60 dof). This latter fit yields an atmospheric temperature (seen at infinity) $T_{\text{eff}} = 61^{+6}_{-12} \text{ eV}$ (90% confidence), a best fit n_{H} consistent with zero and a 90% confidence upper limit of $5 \times 10^{20} \text{ cm}^{-2}$, and flux in the 0.5–3 keV energy band $(1.15 \pm 0.05) \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$. In this model the power-law spectral component accounts for roughly 25% of the total flux in the 0.5–3 keV energy band.

3. CONCLUSIONS

We have identified three new nearby radio MSPs as the counterparts of bright and previously unassociated *Fermi* LAT sources at high Galactic latitude. Our non-detection of young pulsars or MSPs in the more numerous sources searched at low Galactic latitude is likely due to our only moderate sensitivity improvements (typically 2–3 \times) over the best surveys of those regions to date (e.g., Manchester et al. 2001) due to higher sky temperatures resulting from our lower observing frequency. Additionally, the complicated and confused nature of the Galactic plane in γ -rays makes the positive identification of point sources difficult. Several of the bright sources may be blends of other sources or the result of insufficient modeling of the Galactic background. Nonetheless, deeper surveys at frequencies of 1.5–2 GHz of these sources may prove more fruitful in the future.

The new pulsars are very typical radio MSPs in terms of spin period, binary parameters, magnetic field strength, spin-down luminosity, and characteristic age, and their unusual brightness in γ -rays is likely due more to their proximity than to especially energetic emission processes in their magnetospheres. The very high implied γ -ray efficiency for PSR J0614–3329 suggests

that it is likely closer, by up to a factor of two or more, than predicted by the NE2001 model (Cordes & Lazio 2002). The line of sight to PSR J0614–3329 is nearly tangent to the Gum Nebula where NE2001 shows an exceptionally steep DM gradient. Additionally, the pulsar's γ -ray emission is likely not isotropic, but only covers tens of percent of the sky. These large efficiencies in general, though, are consistent with the tens of percent values found by Abdo et al. (2009a) for radio MSPs detected in γ -rays and imply that MSPs are very efficient producers of γ -rays.

We do not have proper motion measurements for pulsars J0614–3329 or J1231–1411 and so their measured spin-down rates are contaminated at some level (likely $\lesssim 10\%$) by the Shklovskii effect (Shklovskii 1970). There is a statistically significant proper motion measurement of $\sim 100 \text{ mas yr}^{-1}$ for J1231–1411, though, which implies a Shklovskii effect at 400 pc larger than the measured spin-down rate for the pulsar. If the proper motion is confirmed at this level, the requirement to have the pulsar intrinsically spinning down gives an upper limit for the pulsar's distance of $\sim 240 \text{ pc}$. Timing observations over the next several years will determine the proper motions and possibly the timing parallaxes for each of the pulsars.

In all three cases we identified X-ray counterparts to the pulsars which substantially aided in the rapid establishment of timing solutions. The three MSPs appear to have fairly typical X-ray properties for radio MSPs (e.g., Bogdanov et al. 2006) with primarily soft thermal-like spectra and X-ray luminosities in the 10^{30} – $10^{31} \text{ erg s}^{-1}$ range, approximately 10^{-4} to 10^{-3} of their γ -ray luminosities.

The radio flux densities of $\sim 1 \text{ mJy}$ near 1 GHz are large enough to make the MSPs potentially useful for a wide variety of timing projects, such as the detection of gravitational waves via long-term pulsar timing (e.g., NANOGrav²⁷), yet they are small enough to explain why earlier large-area surveys for pulsars missed them (e.g., Manchester et al. 1996; Lorimer et al. 2005). In addition, the fact that many of the nearby radio MSPs are being detected in γ -rays and vice versa argues that the sizes of the radio and γ -ray beams are comparable for MSPs (likely within a factor of ~ 2), and that deep radio and γ -ray surveys may allow us to eventually detect a large percentage of the local population of these sources. In the short term, the fact that *Fermi* can point us to nearby radio MSPs is already causing a large increase in the number of known systems, with much less effort than is required by sensitive large-area radio surveys. If most radio MSPs produce γ -rays as these early results seem to indicate, MSPs may contribute to the diffuse isotropic γ -ray background (Faucher-Giguère & Loeb 2010).

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²⁷ <http://nanograv.org>

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