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W mass and Triple Gauge Couplings at Tevatron

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Abstract. The W mass is a crucial parameter in the Standard Model (SM) of particle physics, providing constraints on the mass of the Higgs boson as well as on new physics models via quantum loop corrections. On the other hand, any deviation of the triple gauge boson couplings (TGC) from their values predicted by the SM would be also an indication for new physics. We present recent measurements on W boson mass and searches for anomalous TGC (aTGC) in $W\gamma$, $Z\gamma$, WW , WZ and ZZ at Fermilab Tevatron both by CDF and DØ Collaborations. The CDF Collaboration has measured the W boson mass using data corresponding to $2.2 fb^{-1}$ of integrated luminosity. The measurement, performed using electron and muon decays of W boson, yields a mass of $M_W = 80387 \pm 19$ MeV. The DØ Collaboration has measured $M_W = 80367 \pm 26$ MeV with data corresponding to $4.3 fb^{-1}$ of integrated luminosity in the channel $W \rightarrow e\nu$. The combination with an earlier DØ result, using independent data sample at $1 fb^{-1}$ of integrated luminosity, yields $M_W = 80375 \pm 23$ MeV. The limits on anomalous TGCs parameters are consistent with the SM expectations.

1 W mass measurement at the Tevatron

1.1 Analysis Strategy

At the Tevatron, W bosons are primarily produced in $q\bar{q}$ annihilation, $q\bar{q} \rightarrow W + X$, where X can include QCD radiation that results in measurable hadronic recoil in events. $W \rightarrow l\nu_l$ decays, where $l = e$ or μ , are selected with high purity by the CDF and DØ ($l = e$ only) detectors and used to measure the mass of the W boson.

M_W is determined using three kinematic variables measured in the plane perpendicular to the beam direction: the lepton ($p_T^{e,\mu}$) and neutrino (p_T^ν) transverse momentum distributions and the transverse mass $m_T = \sqrt{2p_T^{e,\mu} p_T^\nu (1 - \cos \Delta\phi)}$, where $\Delta\phi$ is the opening angle between the lepton and neutrino momenta. The magnitude and direction of p_T^ν is inferred from the missing transverse energy (\vec{E}_T).

A parametrized Monte Carlo simulation (FAST MC) is used to predict the shape of the transverse mass distribution as well as the lepton and neutrino p_T distributions as a function of M_W . The W boson mass is extracted by fitting the M_W prediction of the Monte Carlo simulation to the data with a binned maximum-likelihood fit.

The line-shapes prediction depend on the kinematic distributions of the W decay products and detector effects, which are constrained from control samples and theoretical calculations. The kinematic distributions are determined by several effects including internal QED radiation, the intrinsic W boson transverse momentum, and the proton parton distribution functions (PDFs). Electron efficiency, hadronic recoil modelling, calorimeter

response both to electromagnetic shower and hadronic shower and calorimeter fiducial acceptance are among the most important detector effects for both experiments. The parametrized fast simulation enables a study of these effects at a level below 1 part in 10^4 .

1.2 Event Generation and Simulation

Kinematics are simulated using RESBOS [1] which is a next-to-leading order generator including non-perturbative effects at low boson p_T [4]. The PDFs used in RESBOS, are CTEQ6.6 [3]. They affect the W boson mass measurement through their effects on the kinematics of decay charged leptons. The dominant effect from QED radiation to the W boson mass measurement comes from radiation of photons from the final state charged lepton. These processes are simulated by combining the PHOTOS program [2] with RESBOS.

The techniques used by CDF and DØ Collaborations to determine the single most important ingredient, namely the lepton energy scale are very different. The CDF measurement exploits precise measurements of track momenta in the inner tracker, while the DØ measurement relies on a detailed understanding of electron showers in the calorimeter.

At CDF, the high statistics J/Ψ and $\Upsilon(1S) \rightarrow \mu\mu$ quarkonia, are used to set the momentum scale. Fig 1 (left) shows the measurement of the Z -boson mass in the dimuon channel. The Z boson mass measured is $M_Z = 91180 \pm 12_{stat} \pm 10_{syst}$ MeV, in good agreement with the world average (91188 ± 2 MeV).

The electromagnetic calorimeter energy scale is set using the peak of the E/p electron distribution from $W \rightarrow e\nu$

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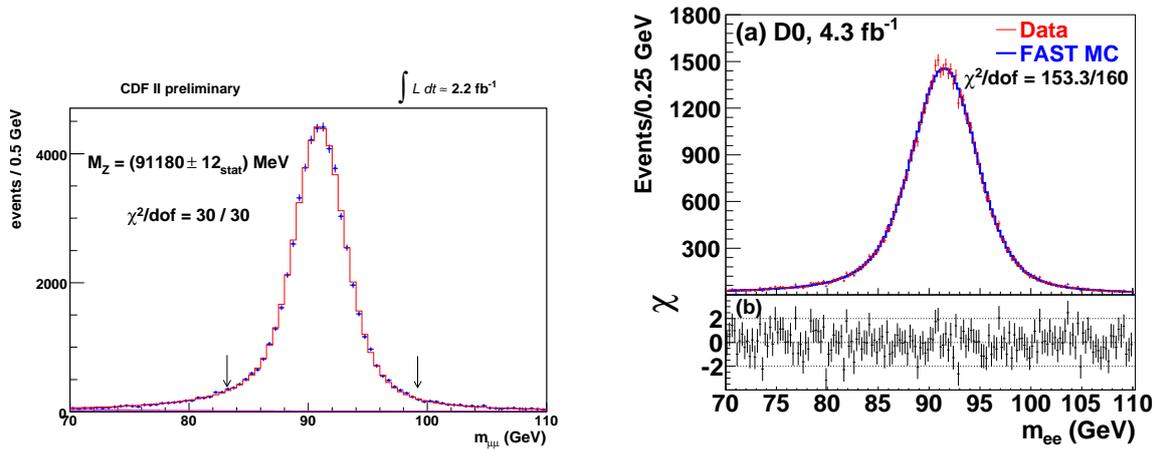


Figure 1. CDF (left): The $Z \rightarrow \mu\mu$ mass fit, showing the data (points) and the best fit simulation template (histogram). The arrows indicate the fitting range. DØ (right):(a) Distribution of the di-electron invariant mass in $Z \rightarrow e^+e^-$ sample. The shape of the distribution of the FAST MC simulation, normalised to the data, is superimposed. (b) Bin-by-bin difference between data and prediction, divided by the statistical uncertainty in the data.

and $Z \rightarrow ee$ samples. The $Z \rightarrow ee$ mass is fitted to cross-check the energy scale and the non-linearity. A value of $M_Z = 91230 \pm 30_{stat} \pm 14_{syst}$ MeV is obtained, consistent with the world average. The tail of the E/p distribution is used to tune the absolute number of radiation lengths in the tracker material.

At DØ the calorimeter energy scale is calibrated using $Z \rightarrow e^+e^-$. The corrections for energy lost in uninstrumented regions are based on a comparison between data and detailed GEANT-based simulation of the DØ detector. The precise measurements of the Z boson mass from LEP [7], is used to determine the absolute energy scale of the calorimeter. The M_W measurement presented here is effectively a measurement of the ratio of W and Z boson masses. Figure 1 (right) shows a comparison of the Z boson mass distributions for data and fast simulation.

The hadronic energy in the event contains the hadronic system recoiling from the W boson, the effects of low energy products from spectator parton collisions and other beam collisions, final state radiation, and energy from the recoil particles that enter the electron selection window. The hadronic response and resolution is calibrated using the mean and width of the η_{imb} distribution in $Z \rightarrow ee$ (and $\mu\mu$ for CDF) events in bins of p_T of the leptons. Here, η_{imb} is defined as the projections of the sum of dielectron (dimuon) transverse momentum and the vector sum of the transverse components of the energies measured in the calorimeter cells excluding those associated with the reconstructed electron.

1.3 Results

Kinematic distributions of background events passing the event selection cuts are included in the templates fits with their estimated normalizations. The backgrounds in the W boson samples are $Z \rightarrow ee$ ($\mu\mu$) in which one e (μ) escapes detection, multijet in which a jet is misidentified as an electron with (\vec{E}_T) arising from misreconstruction, and

Table 1. CDF and DØ : Uncertainties

Uncertainty (MeV)	CDF	DØ
lepton energy scale and resolution	7	17
Hadronic recoil energy scale and resolution	5	6
Backgrounds	3	2
Parton distributions	10	11
QED radiation	4	7
$p_T(W)$ model	5	2
Total systematic uncertainty	15	22
W boson statistics	12	13
Total uncertainty	19	26

$W \rightarrow \tau\nu \rightarrow e(\mu) \nu\nu\nu$. CDF has also evaluated pion and kaon decays in flight and cosmic ray backgrounds for the μ decay channel. Background fractions is about 7% for muons(CDF) and 2% for electrons (CDF and DØ).

Concerning DØ the final M_W uncertainty as well as most of the other uncertainties, are driven by the limited Z statistics (see Table 1). The final M_W uncertainty in CDF relies on precise measurements of track momenta in the inner tracker based on the huge J/Ψ and $\Upsilon(1S)$ statistics (see Table 1). The largest source of uncertainty in the model of W production and decay is the finite precision of the PDFs for both experiments.

The W boson mass is measured by performing a binned maximum-likelihood fit to the lepton p_T^l , neutrino p_T and m_T distributions for each lepton channel (electron channel only for DØ). The result of the fits to the m_T kinematic distribution is shown in Fig. 2 both for CDF (right) and DØ (left).

By combining all six fits and taking into account correlations [8] CDF has published [5] the following result:

$$M_W = 80387 \pm 12_{stat} \pm 15_{syst} = 80387 \pm 19 \text{ MeV.}$$

DØ is using m_T and electron p_T^e in the combination only as the result from \vec{E}_T does not have any significant weight in the combination. The DØ combination yields to:

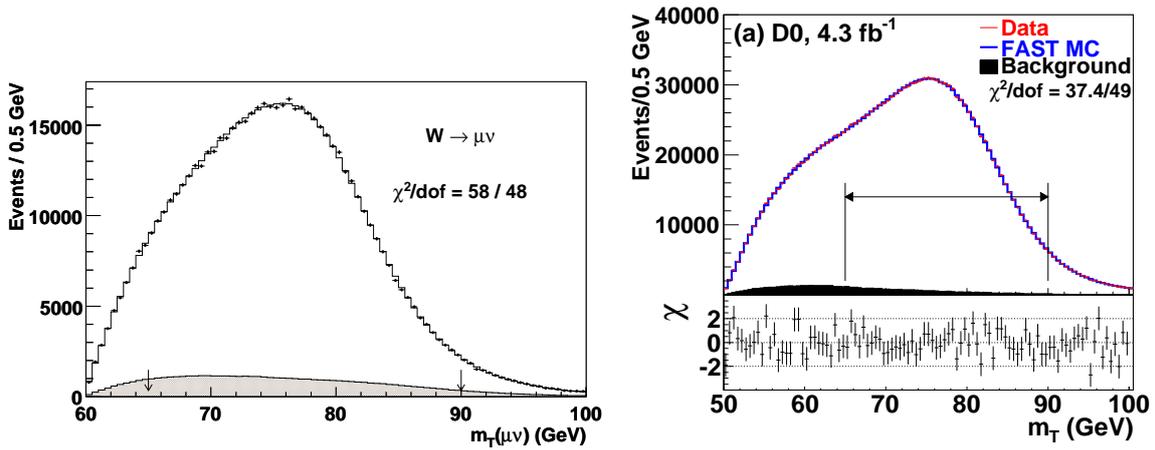


Figure 2. The W transverse mass fit for muon channel in CDF (left). The W transverse mass fit for electron channel in DØ (right). The background contributions to the template are overlaid and arrows indicate the fitting range for both distributions

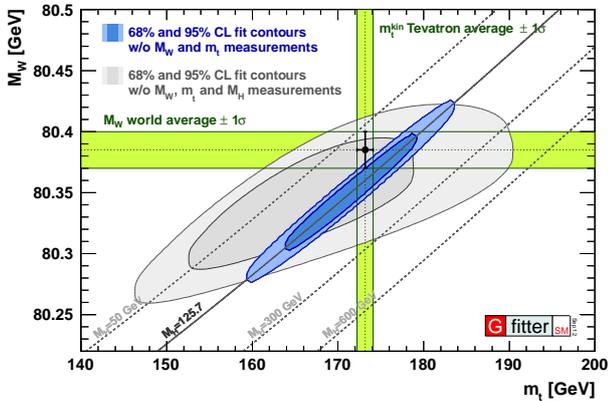


Figure 3. Indirect prediction for the masses of the top quark and the W boson from internal consistency of the SM. This prediction uses the direct measurement of the candidate Higgs boson discovered at the LHC recently. The direct measurements of the W boson mass (the world average discussed in Sec 1.3) and the Tevatron average of direct measurements of the top quark mass [12] are also shown.

$$M_W = 80367 \pm 13_{\text{stat}} \pm 22_{\text{syst}} = 80367 \pm 26 \text{ MeV.}$$

DØ further combines this result [6] with the earlier measurement [9] to obtain the new DØ Run II results:

$$M_W = 80375 \pm 11_{\text{stat}} \pm 20_{\text{syst}} = 80375 \pm 23 \text{ MeV.}$$

The results of the two experiments are of comparable precision, and the measurements are in good agreement. All results from the Tevatron are combined, using again the methods from Ref.[8], to yield the new Tevatron average [10]: $M_W = 80387 \pm 16$ MeV. Combining Tevatron and LEP results [7] yields the new world average: 80385 ± 15 MeV.

1.4 Testing the Standard Model (SM)

Electroweak fit [11] is a powerful tool for the comparison of direct and indirect constraints on the Higgs boson mass.

Moreover, assuming the boson recently discovered at LHC with a mass of 126 GeV is the Higgs boson, electroweak fits can test the validity of the SM. This is shown in Fig. 3 which shows the indirect prediction for the masses of the top quark and the W boson from internal consistency of the SM. This prediction uses the direct measurements of the mass of the candidate Higgs boson. From [11] the indirect constraints prefer a Higgs boson that is consistent within the new boson discovered at LHC at the 1.2σ level, with the mass of the new boson. More precise measurements of M_W are therefore needed to tighten the indirect constraints and make the test of SM even more stringent.

2 Triple Gauge Couplings at Tevatron

2.1 Methods

Measurements of the associated production of two vector bosons (γ , W , Z) are important tests of the electroweak sector of the SM. The production of a diboson final state can occur by particle-antiparticle annihilation (t-channel) or by boson self-interaction (s-channel). Observing triple gauge couplings (TGCs), not permitted in the SM or with different intensity with respect to the SM predictions, would be a sign of new physics.

In the presence of new physics, observables such as production cross-sections, TGCs, and various kinematic distributions, are expected to deviate from their SM predictions. The charged TGCs present in W^+W^- , WZ , and $W\gamma$ production are Δg_1^V , $\Delta \kappa_V$ and λ_V ($V=Z$ or γ) where Δ represents the deviation from SM predictions. In the SM, $\Delta g_1^V = \Delta \kappa_V = \lambda_V = 0$. Using constraints due to gauge invariance, κ_Z and λ_Z can be expressed as a function of the other parameters and the weak mixing angle θ_W as:

$$\kappa_Z = g_Z^1 - (\kappa_\gamma - 1) \tan^2(\theta_W).$$

The neutral TGCs h_i^V ($i=3,4$), studied in $Z\gamma$ production, are not allowed in the SM tree-level, and their values are predicted to be zero.

The p_T or E_T distributions of fully reconstructed candidate bosons or of the charged leptons from boson decays

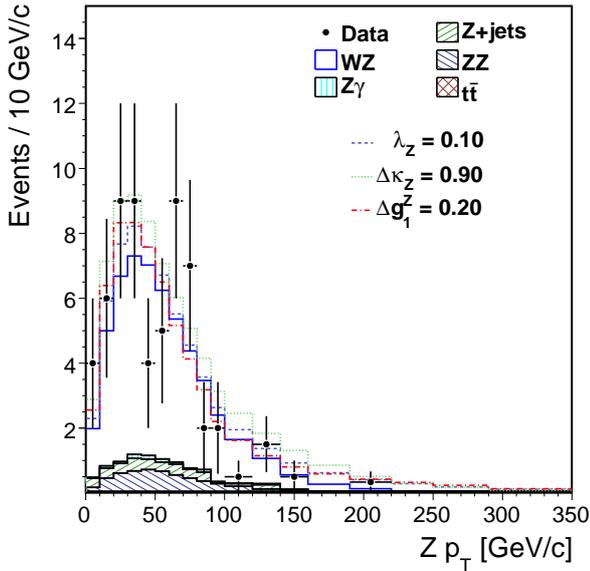


Figure 4. CDF: The $Z p_T$ distribution for data compared to the SM expectation for signal (WZ) and background. Also presented is how the signal expectation would change with anomalous TGCs near the observed limits.

are used to extract limit on aTGC parameters. These distributions are compared with Monte Carlo predictions for different set of coupling constants in order to evaluate 95% C.L. limits. The CDF and $D\phi$ collaborations are producing many competitive measurements in the diboson sector. The most updated results will be presented below.

2.2 Anomalous TGCs limits

The measurement of the production cross section of WZ events in $p\bar{p}$ collisions has been performed by the CDF collaboration [15] with a data sample corresponding to $7.1 fb^{-1}$ of integrated luminosity. The events are reconstructed in final states with three charged leptons and missing transverse energy from a neutrino. The shape and normalization of the p_T spectrum of the Z boson (see Fig. 4) are used to place 95% C.L. limits on anomalous ZWW TGCs. The one-dimensional limits for a cut-off $\Lambda = 2.0$ TeV are:

$$\begin{aligned} -0.39 < \Delta\kappa_Z < 0.90 \\ -0.08 < \Delta\lambda_Z < 0.10 \\ -0.08 < \Delta g_1^Z < 0.20 \end{aligned}$$

the $D\phi$ collaboration has recently published [16] searches of the anomalous γWW and ZWW TGCs from WW and WZ production using lepton plus dijet final states. A combination with results from $W\gamma$ [17], WW [18] and WZ [19] production with leptonic final states has been used to set the most stringent limits at hadron collider to date. The analyzed data correspond to up to $8.6 fb^{-1}$.

The effects of aTGCs is to modify the high boson transverse momentum relative to its SM prediction. Therefore the corresponding p_T^{jj} and p_T^{ll} are used to set 95% C.L. limits on aTGCs (see Fig 5). No deviation from SM is found

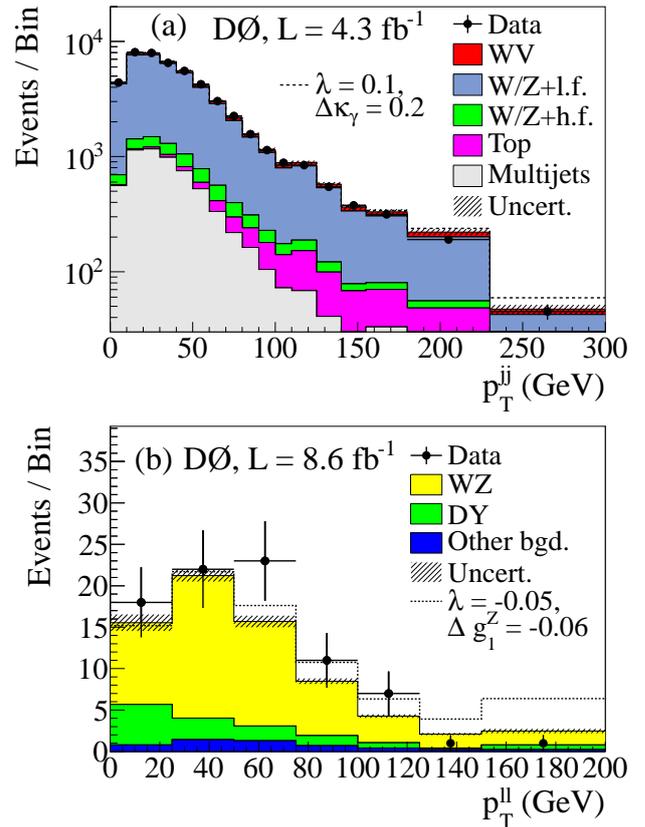


Figure 5. (a) The p_T^{jj} distribution summed over electron and muons channels from $WW + WZ \rightarrow lvjj$ ($l = \mu, e$) production for data and SM MC predictions ("l.f." denotes light-partons and "h.f." denotes heavy-partons). Also shown are expected distributions for an aTGC model with $\Delta\kappa_\gamma = 0.2$ and $\lambda = 0.1$. (b) The p_T^{ll} distribution summed over $eee, e\mu\mu, \mu ee$ and $\mu\mu\mu$ channels from $WZ \rightarrow lvll$ production for data, SM MC prediction and for aTGC model with $\lambda = -0.05$ and $\Delta g_1^Z = -0.06$.

in data. The limits at 95% C.L. from $8.6 fb^{-1} WZ \rightarrow lvll$ are:

$$-0.07 < \lambda < 0.089 \text{ and } -0.055 < \Delta g_1^Z < 0.137$$

while the limits from $4.3 fb^{-1} WW + WZ \rightarrow lvjj$ are:

$$\begin{aligned} -0.27 < \Delta\kappa_\gamma < 0.37 \\ -0.075 < \lambda < 0.080 \\ -0.071 < \Delta g_1^Z < 0.137 \end{aligned}$$

Finally the overall combination gives the following 95% C.L. limits:

$$\begin{aligned} -0.158 < \Delta\kappa_\gamma < 0.255 \\ -0.036 < \lambda < 0.044 \\ -0.034 < \Delta g_1^Z < 0.084 \end{aligned}$$

Based on the combination of all diboson production and decay channels, the most stringent 68% C.L. constraints to date, are set on the W boson magnetic dipole and electric quadrupole moments:

$$\mu_W = 2.012^{+0.035}_{-0.034} \text{ in } e/2M_W \text{ unit}$$

$$q_W = -0.995^{+0.042}_{-0.043} \text{ in } e/M_W^2 \text{ unit.}$$

In beyond-the-SM scenarios, with enhanced $Z\gamma$ couplings, the photon E_T spectrum is modified due to an enhancement in the production of high- E_T photon [20]. The CDF collaboration has used data corresponding to $4.9 fb^{-1}$ in

the $Z \rightarrow \nu\bar{\nu}$ decay channel and $5.1 fb^{-1}$ in the $Z \rightarrow l^+l^-$ ($l=\mu,e$) decay channels [21]. With an energy cut-off of $\Lambda = 1.5$ TeV, the CP-conserving parameters that describe $Z\gamma$ couplings are set to be:

$$|h_3^{\gamma,Z}| < 0.022 \text{ and } |h_4^{\gamma,Z}| < 0.0009$$

These limits are the most stringent limits to date.

The DØ Collaboration has used the transverse momentum spectrum of the photon to place limits on anomalous $ZZ\gamma$ and $Z\gamma\gamma$ couplings [22]. Data corresponding to an integrated luminosity of $6.2 fb^{-1}$ in $p\bar{p} \rightarrow Z\gamma \rightarrow l^+l^-\gamma$ ($l = e, \mu$) production has been analysed. The aTGCs limits are combined with a previous DØ $Z\gamma$ analysis [23] resulting in a total integrated luminosity of $7.2 fb^{-1}$ in the $l\bar{l}\gamma$ decay channel and $3.6 fb^{-1}$ in the $\nu\bar{\nu}\gamma$ decay channel. Applying an energy cut-off of $\Lambda = 1.5$ TeV the limits at 95% C.L. are set to be:

$$|h_3^Z| < 0.026 \text{ and } |h_4^Z| < 0.0013$$

$$|h_3^\gamma| < 0.027 \text{ and } |h_4^\gamma| < 0.0014$$

These limits are comparable to those found in the most recent CDF result reported above.

2.3 Conclusion

A new precision measurement of the W boson mass performed by CDF and DØ Collaborations has improved the precision of the Tevatron average to 16 MeV, which combined with the LEP average slightly reduces the new world average to 15 MeV. These improvements in the experimental precision on M_W lead to precise indirect constraints on the mass of the Higgs boson. The direct measurements of the mass of the Higgs boson candidate discovered at the LHC agree, at 1.2σ level, with indirect constraints. Even if this is a remarkable success for the standard model, more precise measurements of M_W are needed to tighten the indirect constraints and make this test of SM even more stringent. The precision of both CDF and DØ experiments is expected to improve significantly once the full data set (almost $10 fb^{-1}$) has been analysed.

The CDF and DØ Collaborations have presented recent results of diboson production using the data collected at the Tevatron between 2002-2011. The limits on anomalous TGCs parameters are competitive with previous results from the LEP and first results from the LHC. They are consistent with the SM expectations. CDF and DØ can anticipate 2 to 10 times more data (depending on the channels), to cope with the full integrated luminosity ($10 fb^{-1}$) at Tevatron. This will result in an increase of the aTGCs sensitivities by a factor 3 to 5 after a combination of the CDF and DØ results.

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