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Leptons in Near Earth Orbit

The AMS Collaboration

Abstract

The lepton spectra in the kinetic energy ranges 0.2 to 40 GeV for e^- and 0.2 to 3 GeV for e^+ were measured by the Alpha Magnetic Spectrometer (AMS) during space shuttle flight STS-91 at altitudes near 380 km. From the origin of the leptons two distinct spectra were observed: a higher energy spectrum and a substantial second spectrum with positrons much more abundant than electrons. Tracing leptons from the second spectra shows that most of these leptons travel for an extended period of time in the geomagnetic field and that the e^+ and e^- originate from two complementary geographic regions.

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Introduction

The current understanding of the high energy lepton (e^\pm) spectra in cosmic rays is that they are dominated by an electron component. High energy electrons are believed to originate from primary acceleration sites, specifically from supernova explosions. High energy electron–positron pairs are thought to be produced from the collisions of cosmic ray hadrons and gamma rays with interstellar gas. Taken together, the expected positron to electron ratio in cosmic rays arriving at Earth is roughly 10% and it decreases with energy. This picture is based on the experimental data collected over 35 years [1,2] by balloon experiments as well as phenomenological model descriptions developed over the same period [3]. These experiments were performed at altitudes of 30–40 km. Balloon experiments have made important contributions to the understanding of primary cosmic ray spectra and the behavior of atmospheric secondary particles in the upper layer of the atmosphere.

A few pioneering satellite experiments [4] have reported data on low energy electrons and positrons trapped in the geomagnetic field. The satellite based detectors used so far, *i.e.* before this experiment, have not been sensitive enough to systematically study the electron and positron spectra over a broad energy range and their dependence on position and angle.

The electron spectrum observed near Earth shows a low energy drop off due to the geomagnetic cutoff. Previous measurements above the cutoff indicate that the spectrum falls off according to a power law.

The Alpha Magnetic Spectrometer (AMS) [5] is a high energy physics experiment scheduled for installation on the International Space Station. In preparation for this mission, AMS flew a precursor mission on board the space shuttle Discovery during flight STS–91 in June 1998. In this report we use the data collected to study the spectra of electrons and positrons in cosmic rays over the respective kinetic energy ranges of 0.2 to 40 GeV and 0.2 to 3 GeV, the latter range being limited by the proton background.

The large acceptance of AMS and high statistics ($\sim 10^5$) enable us to study the variation of the spectra with position and angle both above and below the geomagnetic cutoff. The accurate momentum resolution, precise trajectory reconstruction and good particle identification of AMS allow an investigation into the origin of particles below cutoff by tracking them in the geomagnetic field.

The AMS detector

The major elements of AMS as flown on STS–91 were a permanent magnet, a tracker, time of flight hodoscopes, a Cerenkov counter and anti-coincidence counters [6,7]. The permanent magnet had the shape of a cylindrical shell with inner diameter 1.1 m, length 0.8 m. It provided a central dipole field of 0.14 Tesla across the magnet bore and an analyzing power, BL^2 , of 0.14 Tm^2 parallel to the magnet, or z –, axis. The six layers of double sided silicon tracker were arrayed transverse to the magnet axis. The outer layers were just outside the magnet bore. The tracker measured the trajectory of relativistic singly charged particles with an accuracy of 20 microns in the bending coordinate and 33 microns in the non-bending coordinate, as well as providing multiple measurements of the energy loss. The time of flight system had two planes at each end of the magnet, covering the outer tracker layers. Together the four planes measured singly charged particle transit times with an accuracy of 120 psec and also yielded multiple energy loss measurements. Two layers of Aerogel threshold Cerenkov counter with an index of refraction $n = 1.035$ were used to make independent velocity measurements allowing the discrimination of lower energy hadrons from electrons and positrons. A layer of anti-coincidence scintillation counters lined the inner surface of the magnet. Low energy particles were absorbed by thin carbon fiber shields. In flight the AMS positive z –axis pointed out of the shuttle payload bay.

For this study the acceptance was restricted to events with an incident angle within 25° of the positive z -axis of AMS and data from four periods are included. In the first period the z -axis was pointing within 1° of the zenith. Events from this period are referred to as “downward” going. In the second period the z -axis pointing was within 1° of the nadir. Data from this period are referred to as “upward” going. In the third and fourth periods the AMS z -axis was pointing within 20° and 45° of the zenith. The orbital inclination was 51.7° and the geodetic altitude during these periods ranged from 350 to 390 km. Data taken while passing through or near the South Atlantic Anomaly were excluded from this analysis.

The response of the detector was simulated using the AMS detector simulation program, which is based on the GEANT package [8]. The effects of energy loss, multiple scattering, interactions, decays and the measured detector efficiency and resolution were included.

After the flight the AMS detector was extensively calibrated at two accelerators: at GSI, Darmstadt, with helium and carbon beams at 600 incident angles and locations and 10^7 events, and at the CERN proton-synchrotron in the energy region of 2 to 14 GeV, with 1200 incident angles and locations and 10^8 events. This ensured that the performance of the detector and the analysis procedure were thoroughly understood.

Analysis

Event reconstruction, analysis and spectrum unfolding are detailed in [9]. Electron candidates were specifically selected by requiring the measured particle charge to be -1 and the particle velocity to be compatible with the speed of light. Backgrounds arose from protons with wrongly measured momentum and secondary pions produced in the detector materials. The two most important cuts used to remove these backgrounds were on the χ^2 value obtained in fitting the particle trajectory, which removed tracks with large single or multiple scattering, and on the number of hits near the reconstructed trajectory in both the tracker and time of flight scintillators.

After the above cuts were applied, the overall probability of a proton event to be accepted as an electron, estimated from Monte Carlo simulations and confirmed in the CERN test beam, was $O(10^{-4})$ with an electron selection efficiency of 75%. To further reduce the pion background only events whose track passed through the active Cerenkov counter area and, therefore, had an independent velocity measurement were accepted.

Positron candidates were selected by requiring the charge to be $+1$ and, as for electrons, the velocity be compatible with the speed of light and track quality cuts. In contrast to electrons, the main background for the positron sample came from proton events with poorly measured velocity. The rejection power against this background decreased rapidly with increasing proton momentum, therefore tighter quality cuts on the velocity measurements were applied. Above $1 \text{ GeV}/c$ protons were rejected by requiring two independent velocity measurements from the two separate Cerenkov counter layers to be compatible with the velocity of a positron. Lower energy protons were rejected by requiring the energy loss measurements in four layers of time of flight counters and six double layers of silicon tracker to be compatible with a positron. These cuts yielded an additional background rejection factor of 5 at the expense of lower positron selection efficiency. Table 1 summarizes the estimated efficiencies.

A convolution of the background rejection function with the measured proton spectra provided an energy dependent background estimation. Fig. 1 shows the measured electron and positron spectra together with the estimated background for the geomagnetic polar regions, where the background conditions were most severe.

Cut	Efficiency (%)
Tracking Quality Cuts	75 ± 3
Common e^\pm Velocity Cuts	52 ± 1
Additional e^+ Velocity Cuts	72 ± 1.5
Total electrons	39 ± 1.7
Total positrons	28 ± 1.3

Table 1: Percentage e^\pm selection efficiencies and uncertainties.

The acceptance was determined as a function of particle momentum and direction. The average acceptance was found to rise from about $0.01 \text{ m}^2 \text{ sr}$ at 0.15 GeV and level off at $0.1 \text{ m}^2 \text{ sr}$ above 0.7 GeV with a systematic uncertainty of 5% [9]. The incident differential spectrum was obtained from the measured spectrum by using an unfolding method based on Bayes' theorem [10] with resolution functions obtained from the simulation. These functions were confirmed at several energy points with calibration measurements in the CERN proton beams.

Results and interpretation

Fig. 2 presents the downward lepton spectra integrated over incident angles within 25° of the AMS z -axis, which was within 1° of the zenith. In Fig. 3 these spectra are compared with the spectra measured with upward going leptons. The measurements have been binned according to the absolute value of the corrected geomagnetic latitude [11], Θ_M (radians), at which they were detected. The effect of the geomagnetic cutoff and the decrease in this cutoff with increasing Θ_M is particularly visible in the downward electron spectra. The spectra above and below cutoff differ. To understand this difference the trajectory of electrons and positrons were traced [12] back from their measured incident angle, location and momentum, through the geomagnetic field [13]. This was continued until the trajectory was traced to outside the Earth's magnetosphere or until it crossed the top of the atmosphere at an altitude of 40 km. In a refinement from [9], the spectra from particles which were traced to originate far away from Earth are classified as "primary" and those from particles which originate in the atmosphere as "second" spectra. In practice particles below the cutoff are from the second spectra, however this classification provides a cleaner separation in the transition region.

I. Properties of the primary lepton spectra

Fig. 4a shows the primary lepton spectra. The spectra are in reasonable agreement with previous measurements [2]. Fig. 4b shows the energy dependence of the positron fraction, which exhibits the predominance of electrons over positrons in primary cosmic rays.

II. Properties of the second lepton spectra

As shown in Figs. 2 and 3, substantial second lepton spectra are observed for downward and upward going leptons at all geomagnetic latitudes below the geomagnetic cutoff. These spectra have the following properties:

- (i) The second lepton spectra of Fig. 2 exhibit similar qualitative behavior to the proton spectra [9].

- (ii) At polar latitudes the downward second spectrum of electrons is gradually obscured by the primary spectrum, whereas the second spectrum of upward going electrons is clearly observed (see Figs. 2 and 3).
- (iii) For both electrons and positrons the upward and downward fluxes are nearly identical (see Fig. 3).
- (iv) As seen from Fig. 5 the lepton fluxes reach a maximum at the geomagnetic equator. With increasing latitude the positron flux drops off faster than the electron flux.

In addition to the backward tracing mentioned above the leptons were also traced forward until their trajectory would have either escaped or crossed the top of the atmosphere, the location of which was taken as the particle sink. The results show that all second spectrum particles eventually re-enter the atmosphere. Defining the flight time as the sum of forward and backward tracing times, that is the interval between origin and sink, Fig. 6 shows the distribution of flight time versus energy for electrons and positrons. Both e^+ and e^- exhibit two distinct types of trajectories:

- The horizontal bands with flight times < 0.2 sec, defined as “short-lived”.
- The diagonal bands with flight times ≥ 0.2 sec defined as “long-lived”.

For $\Theta_M < 0.3$, most (75% of e^+ , 65% of e^-) leptons are long-lived.

Distinct properties of the second spectra for short-lived leptons

The trajectory tracing shows that leptons travel in cycles across the equator where the trajectories reach maximal altitude and they are reflected at the lowest points at the mid and polar latitudes. For short-lived leptons:

- From Fig. 6 one sees that the flight time is independent of lepton energy.
- The point of origin shows no longitude dependence. They do not originate from near to the geomagnetic equator, $\Theta_M < 0.4$ (see Fig. 7a,b).
- The particle flux is independent of the shuttle attitude and is approximately isotropic (see Fig. 7c,d,e).

Distinct properties of the second spectra for long-lived leptons

- As shown in Fig. 8 long-lived e^- and e^+ originate from well defined, complementary geographic regions. Tracing also shows that the regions of origin for positrons coincide with regions of sink for electrons and vice versa.
- Fig. 9 shows the strongly peaked distributions of the point of origin of the long-lived leptons in geomagnetic coordinates. Within the regions indicated the distributions are strongly peaked and the two diagonal bands (A , B) seen in Fig. 6 for the long-lived leptons correspond to the two regions of origin (A , B) marked in Figs. 8 and 9.
- The long-lived leptons are reflected across the equator hundreds of times. The number of cycles they can make before being absorbed in the atmosphere decreases with their energy.

- As shown in Fig. 8c,d,e, the long-lived lepton flux reaches a maximum in the equatorial region where they are produced and absorbed.
- At zenith shuttle orientation, 99% of the long-lived leptons are actually detected at $\Theta_M < 0.4$, indicating a strongly anisotropic angular distribution.

We note that the behaviour of protons and positrons is very similar (see [9]).

Lepton charge ratio

An interesting feature of the observed second lepton spectra is the predominance of positrons over electrons. In table 2 the e^+/e^- ratios grouped according to magnetic latitude region and shuttle attitude (0° , 20° , 45° , 180°) are given separately for long-lived and short-lived leptons. As seen from table 2

e^+/e^-	Long-lived (flight time ≥ 0.2 seconds)				
Attitude	$0.0 < \Theta_M < 0.2$	$0.2 < \Theta_M < 0.4$	$0.4 < \Theta_M < 0.6$	$0.6 < \Theta_M < 0.8$	$0.8 < \Theta_M < 1.0$
0°	4.27 ± 0.17	3.26 ± 0.37	1.65 ± 1.24		
20°	4.15 ± 0.39	2.75 ± 0.45	2.92 ± 1.00	1.05 ± 0.69	1.46 ± 0.42
45°	4.36 ± 0.40	3.41 ± 0.30	3.81 ± 0.33	2.27 ± 0.18	1.28 ± 0.16
180°	4.27 ± 0.25	4.25 ± 0.65			
e^+/e^-	Short-lived (flight time < 0.2 seconds)				
Attitude	$0.0 < \Theta_M < 0.2$	$0.2 < \Theta_M < 0.4$	$0.4 < \Theta_M < 0.6$	$0.6 < \Theta_M < 0.8$	$0.8 < \Theta_M < 1.0$
0°	3.08 ± 0.35	2.43 ± 0.19	1.35 ± 0.11	1.10 ± 0.11	0.83 ± 0.10
20°	2.83 ± 0.67	2.23 ± 0.37	1.95 ± 0.28	1.48 ± 0.22	0.94 ± 0.18
45°	3.22 ± 0.44	2.18 ± 0.32	2.01 ± 0.32	1.08 ± 0.12	0.93 ± 0.19
180°	4.84 ± 0.81	2.79 ± 0.28	1.45 ± 0.18	1.17 ± 0.21	0.68 ± 0.27

Table 2: Lepton charge ratio versus magnetic latitude for the shuttle attitudes 0° , 20° , 45° and 180° for long-lived and short-lived particles.

the ratios:

- Depend at most weakly on the shuttle orientation.
- The ratios for short- and long-lived leptons behave differently. For short-lived leptons the e^+/e^- ratio is maximal at the magnetic equator where it reaches a value of ~ 3 whereas for long-lived leptons the ratio is higher, ≥ 4 at the magnetic equator, and less dependent on latitude.
- The energy dependence of the e^+/e^- ratio for 0° attitude and $\Theta_M < 0.3$ is shown in Fig. 10. As seen, short-lived and long-lived leptons behave differently. For short-lived leptons the ratio does not depend on the particle energy in the range 0.2 to 3 GeV but for long-lived leptons the ratio does depend on the lepton energy, reaching a maximum value of ~ 5 .

The combined (short- and long-lived, all attitudes) dependence on Θ_M of the ratio for all second spectra particles is shown in Fig. 11.

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References

- [1] R.R. Daniel and S.A. Stephens, *Phys. Rev. Lett.* **15** (1965) 769; C.J. Bland *et al.*, *Phys. Rev. Lett.* **17** (1966) 813; S.D. Verma, *J. Geophys. Res.* **72** (1967) 915; C.J. Bland *et al.*, *Nouvo Cim.* **LV B** (1968) 451; B. Agrinier *et al.*, *Nouvo Cim.Lett.* **1** (1969) 54; J.L. Faselow *et al.*, *ApJ* **158** (1969) 771; J. Daugherty *et al.*, *ApJ* **198** (1975) 493; A. Buffington *et al.*, *ApJ* **199** (1975) 669; R. Hartman and C. Pellerin, *ApJ* **204** (1976) 927; K.K. Tang, *ApJ* **278** (1984) 881; R.L. Golden *et al.*, *ApJ* **287** (1985) 662; D. Muller and K. Tang, *ApJ* **312** (1987) 183; G. Barbiellini *et al.*, *Astronomy and Astrophysics* **309** (1996) L15; R.L. Golden *et al.*, *ApJ* **457** (1996) L103; S.W. Barwick *et al.*, *ApJ* **482** (1997) L191; S.W. Barwick *et al.*, *J. Geophys. Res.* **103** (1998) 4817; S. Torii *et al.*, *Proc. 26th ICRC* **3** (1999) 53; M. Boezio *et al.*, *Proc. 26th ICRC* **3** (1999) 57; S. Coutu *et al.*, *Astropart.Phys.* **11** (1999) 429.
- [2] R.L. Golden *et al.*, *ApJ* **436** (1994) 769; S.W. Barwick *et al.*, *ApJ* **498** (1998) 779; M.A. DuVernois *et al.*, *Proc. 26th ICRC* **3** (1999) 49.
- [3] O.A. Bogdanova *et al.*, *15 ICRC Plovdiv* **3** (1977) 176; M. Giler *et al.*, *J.Phys.A:Math.Gen.* **10** (1977) 843; R.J. Protheroe, *ApJ* **254** (1982) 391; W.R. Webber, *20 ICRC* **2** (1987) 80; I.V. Moskalenko and A.W. Strong, *ApJ* **493** (1998) 694.
- [4] J.B. Cladis *et al.*, *J. Geophys. Res.* **66** (1961) 2297; L.V. Kurnosova *et al.*, *15 ICRC Plovdiv* **4** (1977) 185; R.N. Basilova *et al.*, *16 ICRC Kyoto* **3** (1979) 150; N.L. Grigorov *et al.*, *Dokl.Akad.Nauk SSSR* **282** (1985) 81; Yu.E. Efimov *et al.*, *Czechoslovak Journ. of Phys.* **35** (1985) 1371; S.A. Voronov *et al.*, *Izv.Vysshikh Uchebn.Zavedenii,Fizika* **9** (1986) 19–24; S.A. Voronov *et al.*, *Geomagnetism and Aeronomy* **27** (1987) 424; A.F. Iyidin *et al.*, *Geomagnetism and Aeronomy* **28** (1988) 103; S.V. Koldashov *et al.*, *24 ICRC Roma* **4** (1995) 993; A.M. Galper *et al.*, *25 ICRC Durban* **4** (1997) 333.
- [5] S. Ahlen *et al.*, *Nucl. Inst. Meth.* **A350** (1994) 351.
- [6] G.M. Viertel and M. Capell, *Nucl. Inst. Meth.* **A419** (1998) 295–299.
- [7] AMS Collaboration, J. Alcaraz *et al.*, *Phys. Lett.* **B461** (1999) 387–396.
- [8] R. Brun *et al.*, *GEANT 3*, CERN DD/EE/84-1 (Revised, 1987); P.A. Aamio *et al.*, *FLUKA Users Guide*, CERN TIS-RP-190 (1990).
- [9] AMS Collaboration, J. Alcaraz *et al.*, *Phys. Lett.* **B472** (2000) 215–226.

- [10] A. Kondor, Nucl. Inst. Meth. **216** (1983) 177; G. D’Agostini, Nucl. Inst. Meth. **A362** (1995) 487.
- [11] A. Brekke, Physics of the Upper Polar Atmosphere, pp. 127–145, (Wiley, 1997).
- [12] Y.L. Chuang *et al.*, Chinese Journal of Physics, in preparation (2000); N. Zographos, ETHZ-IPP 99-04 (1999).
- [13] N.A. Tsyganenko and A.V. Usmanov, Planet. Space Sci. **30** (1982) 985–998; N.A. Tsyganenko *et al.*, Software for Computations of Geomagnetic Field and Related Coordinate Systems, Soviet Geophysical Committee, Special Report (1987); N.A. Tsyganenko, Planet. Space Sci. **35** (1987) 1347–1358; N.A. Tsyganenko, Planet. Space Sci. **37** (1989) 5–20; N.A. Tsyganenko, J. Geophys. Res. **100** (1995) 5599–5612; N.A. Tsyganenko and D.P. Stern, J. Geophys. Res. **101** (1996) 27187–27198; R.L. Langel, Chairman, IAGA Div. V working group 8, J. Geomag. Geoelectr. **47** (1995) 1251–1261; G. Gustafsson, N.E. Papitashvili and V.O. Papitashvili, J. Atmos. Terr. Phys. **54** (1992) 1609–1631.

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- ⁹ Also supported by the Italian Space Agency.

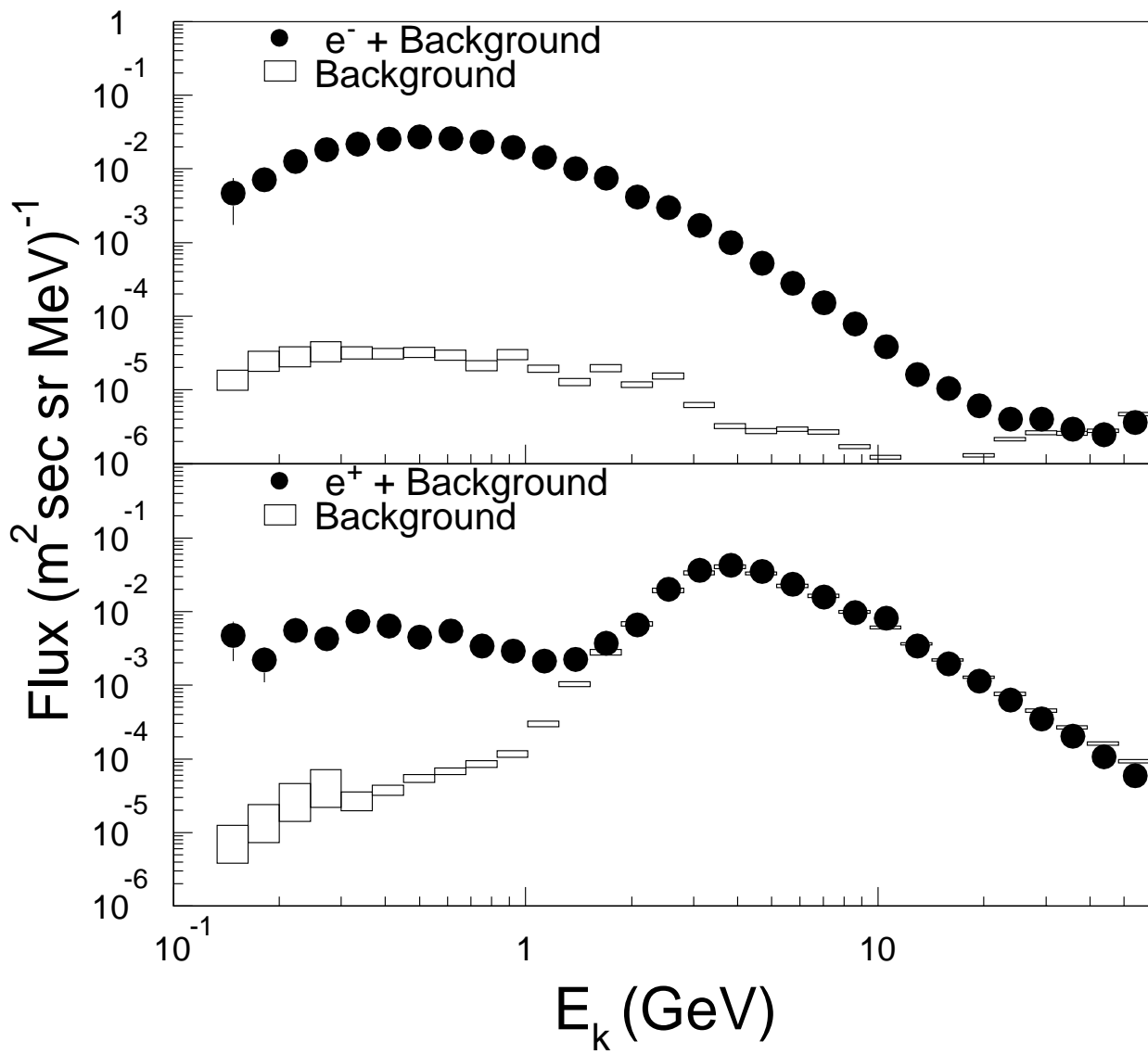


Figure 1: The primary e^\pm fluxes and background in the geomagnetic polar region ($\Theta_M > 0.9$).

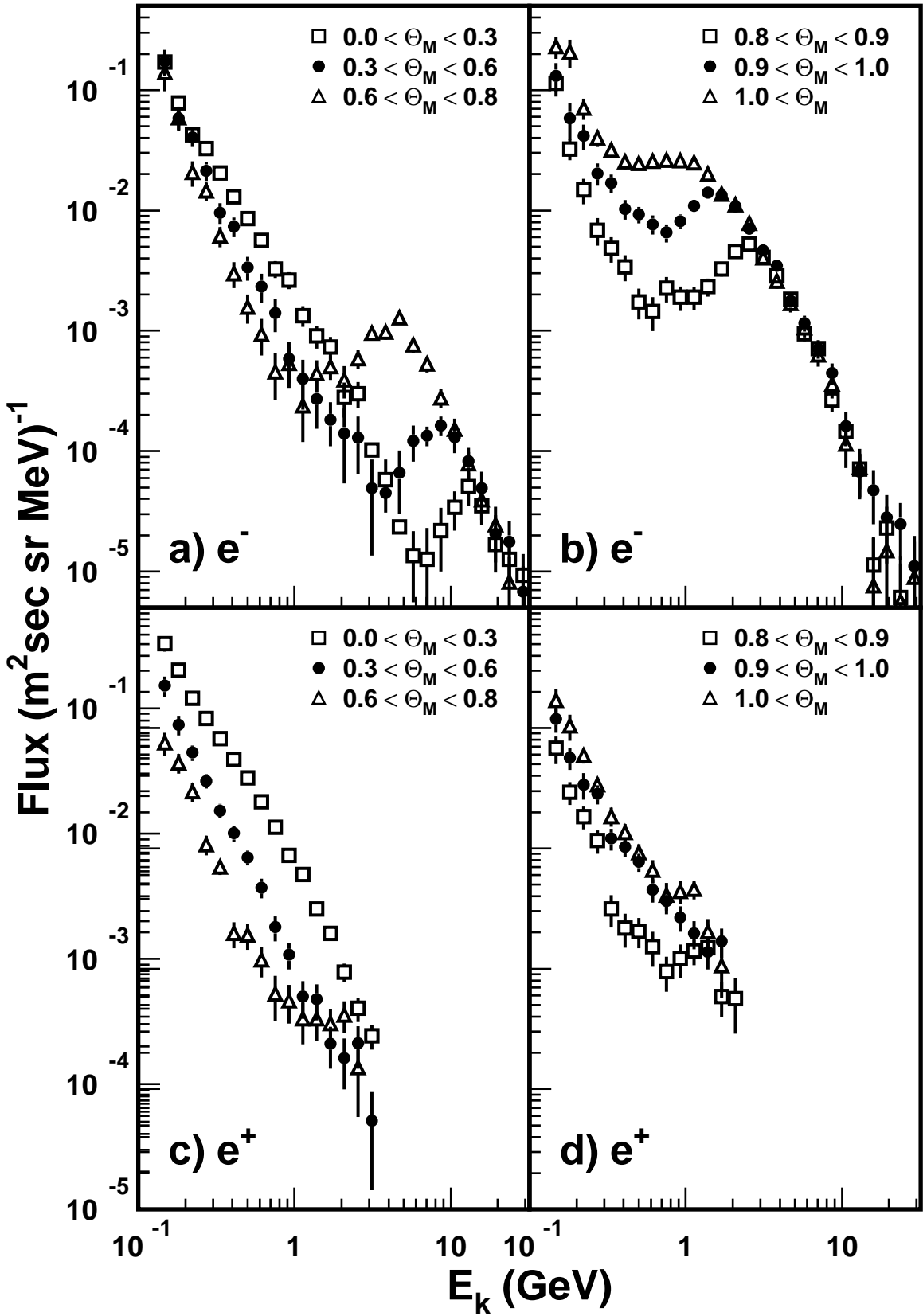


Figure 2: (a,b) Flux spectra for downward going electrons and (c,d) positrons, separated according to the geomagnetic latitude at which they were detected.

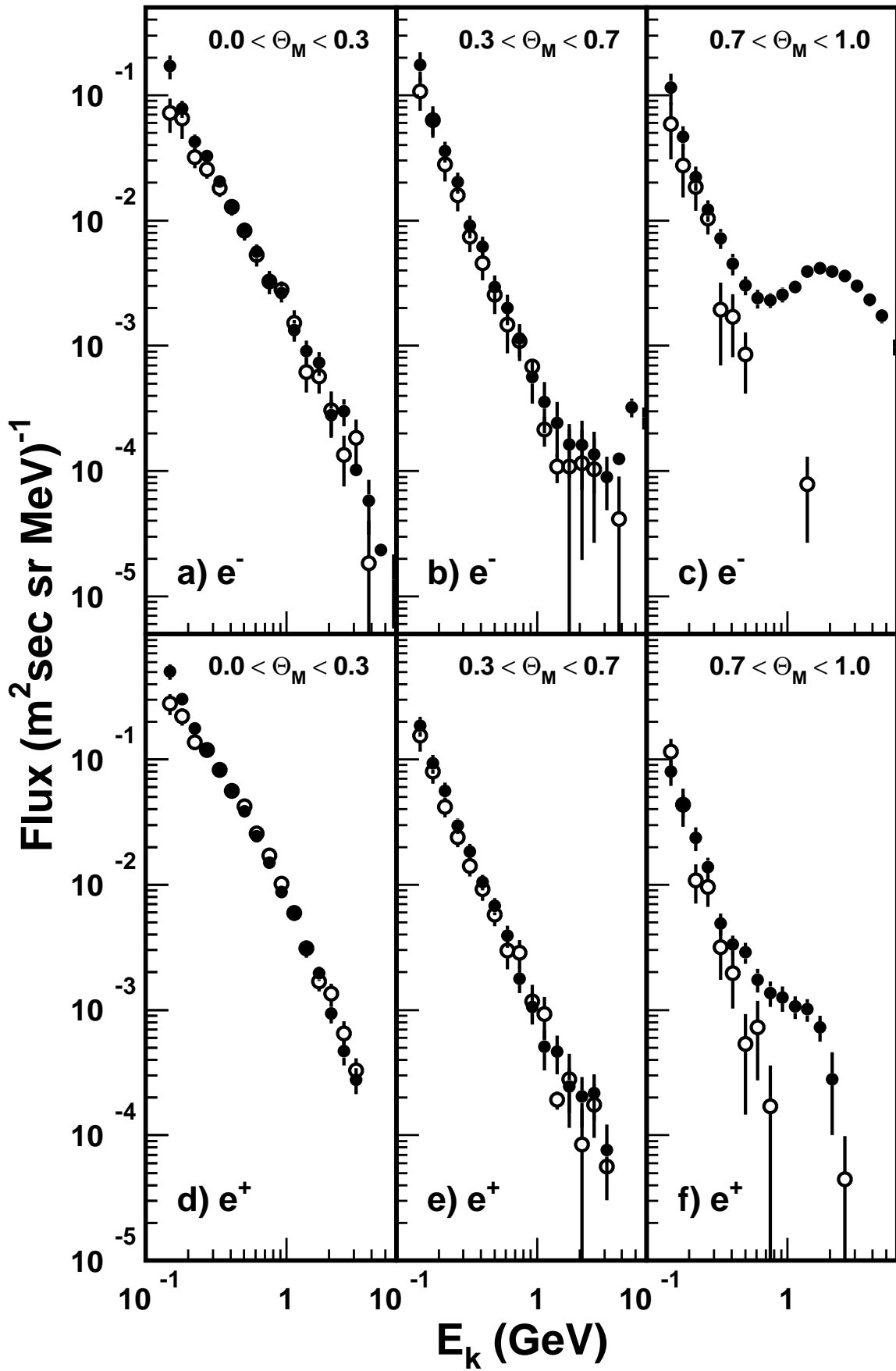


Figure 3: (a,b,c) Flux spectra for downward (full circles) and upward (open circles) going electrons and (d,e,f) positrons, separated according to the geomagnetic latitude, Θ_M , at which they were detected.

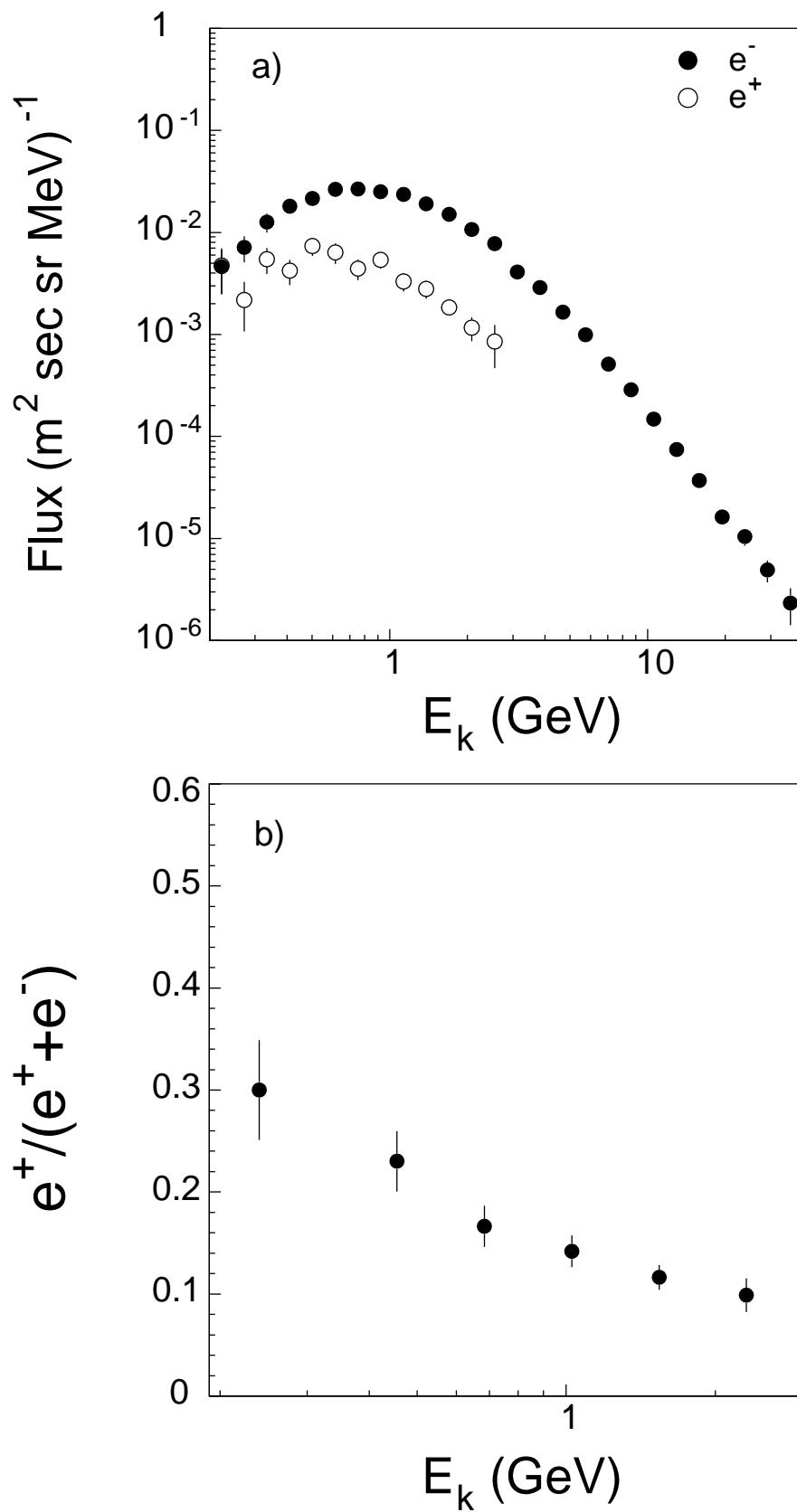


Figure 4: (a) Flux spectra for primary leptons. Particle direction within 25° of zenith. (b) Positron fraction for primary leptons versus energy.

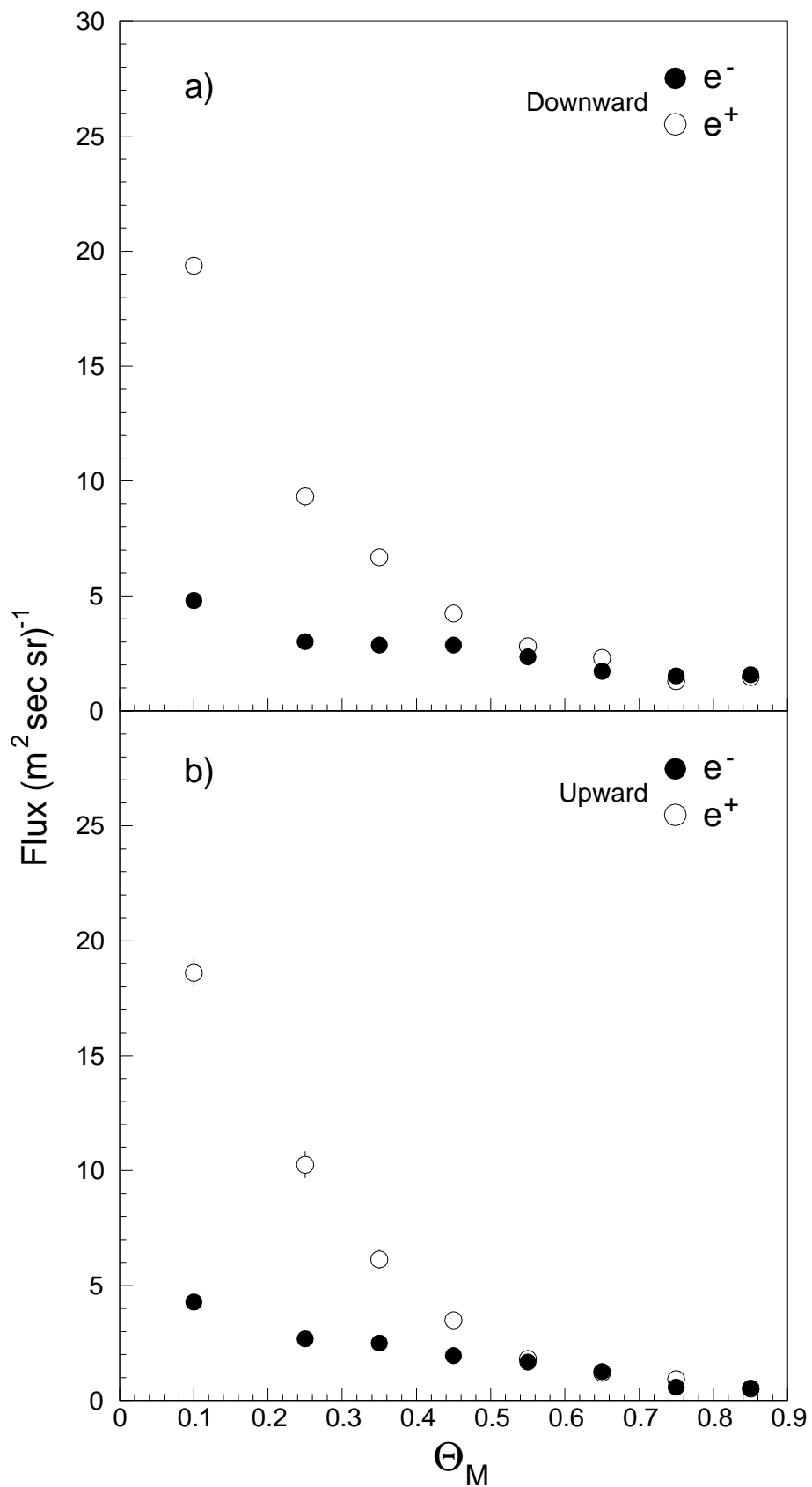


Figure 5: Properties of second lepton spectra flux: (a) downward and (b) upward going electrons and positrons as functions of the geomagnetic latitude, Θ_M , at which they were detected integrated over the range 0.2-2.5 GeV.

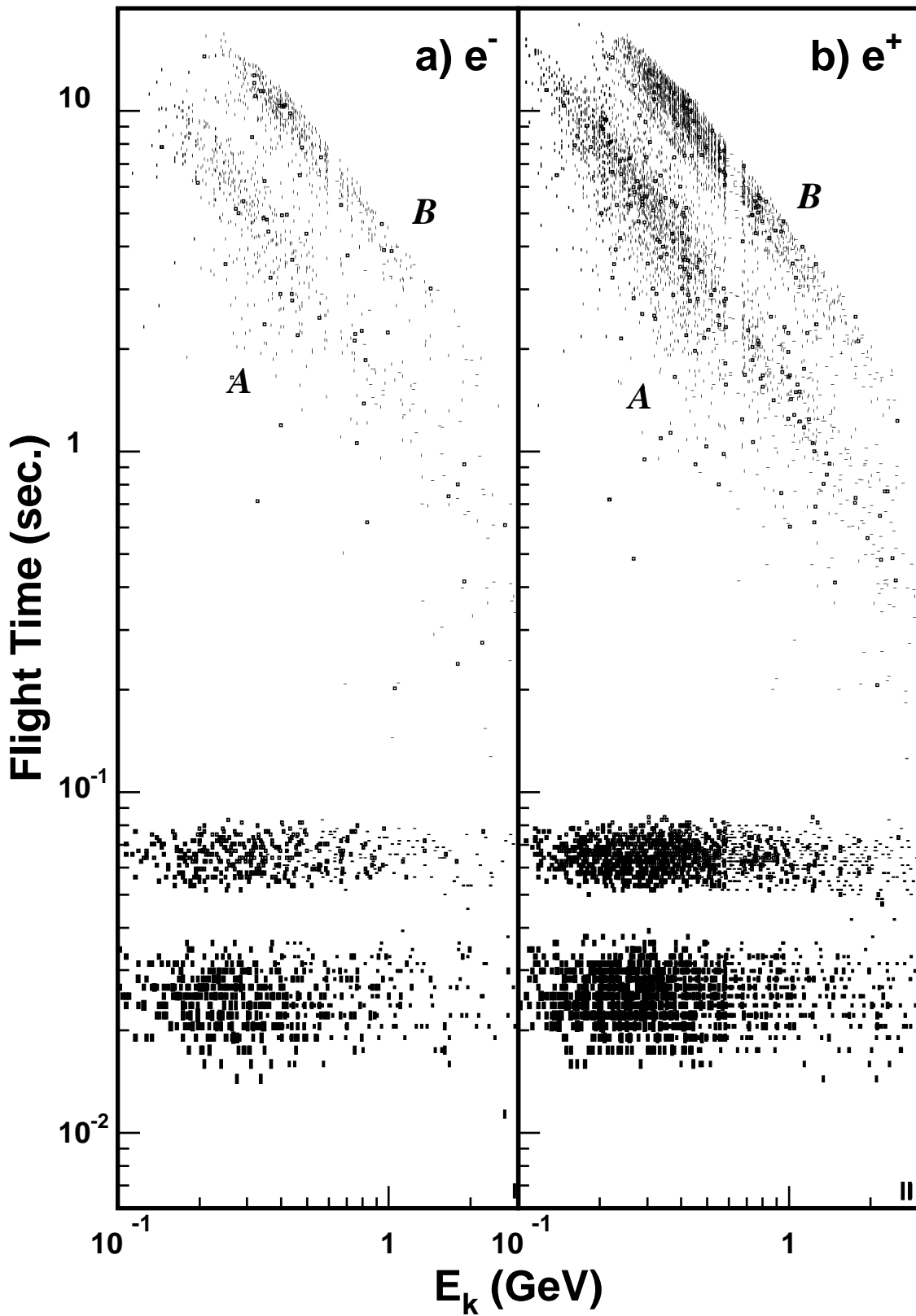


Figure 6: The flight time versus energy from the tracing of leptons detected in the region $\Theta_M < 0.3$. From the flight time distribution there are two distinct types of trajectories: For “short-lived”, flight times < 0.2 sec, the flight time is independent of lepton energy. For “long-lived”, flight times ≥ 0.2 sec, there are two bands A and B. In both A and B the flight time depends on energy: it decreases with increasing energy.

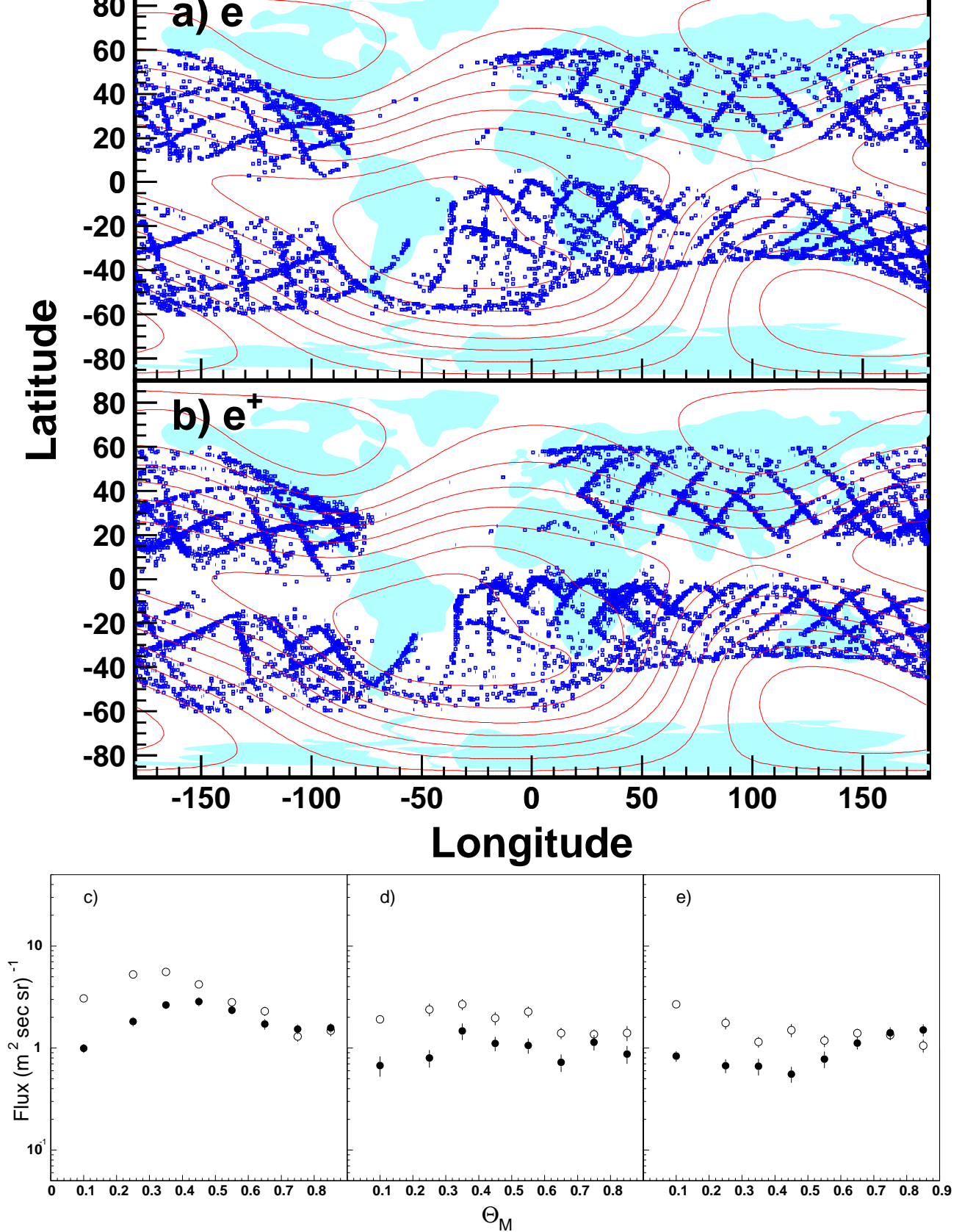


Figure 7: Properties of short-lived second spectra leptons (< 3 GeV): (a) The geographic origin of electrons and (b) positrons. Note that the point of origin shows no longitudinal dependence and that the short-lived leptons do not originate from the region $\Theta_M < 0.4$. The lines indicate the geomagnetic field contours at 380 km. (c) The e^- (full circles) and e^+ (open circles) fluxes integrated over the range 0.2-2.5 GeV as a function of magnetic latitude for zenith, (d) 20° and (e) 45° shuttle attitude.

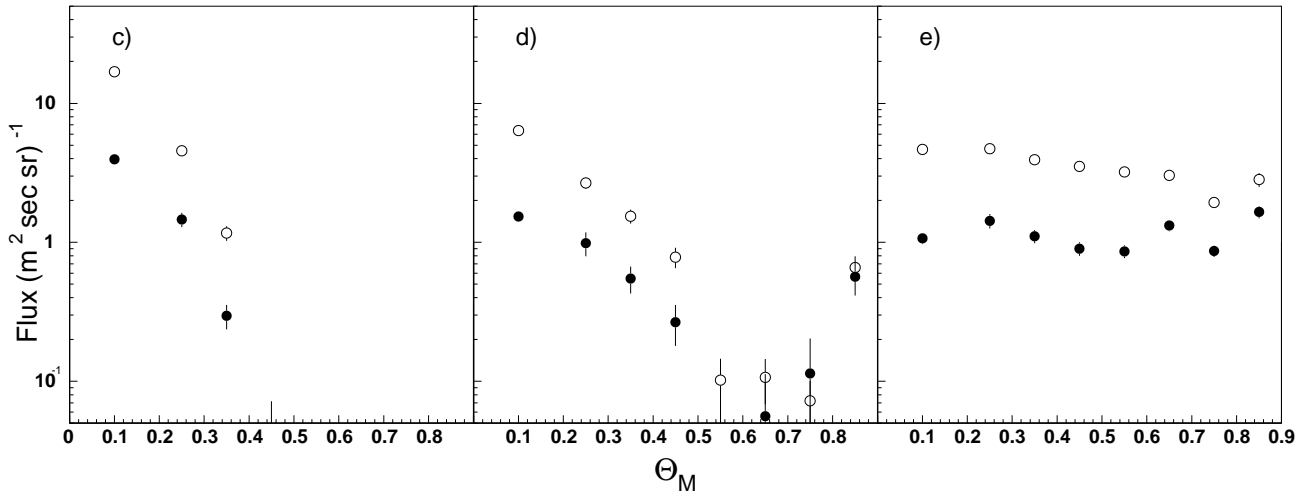
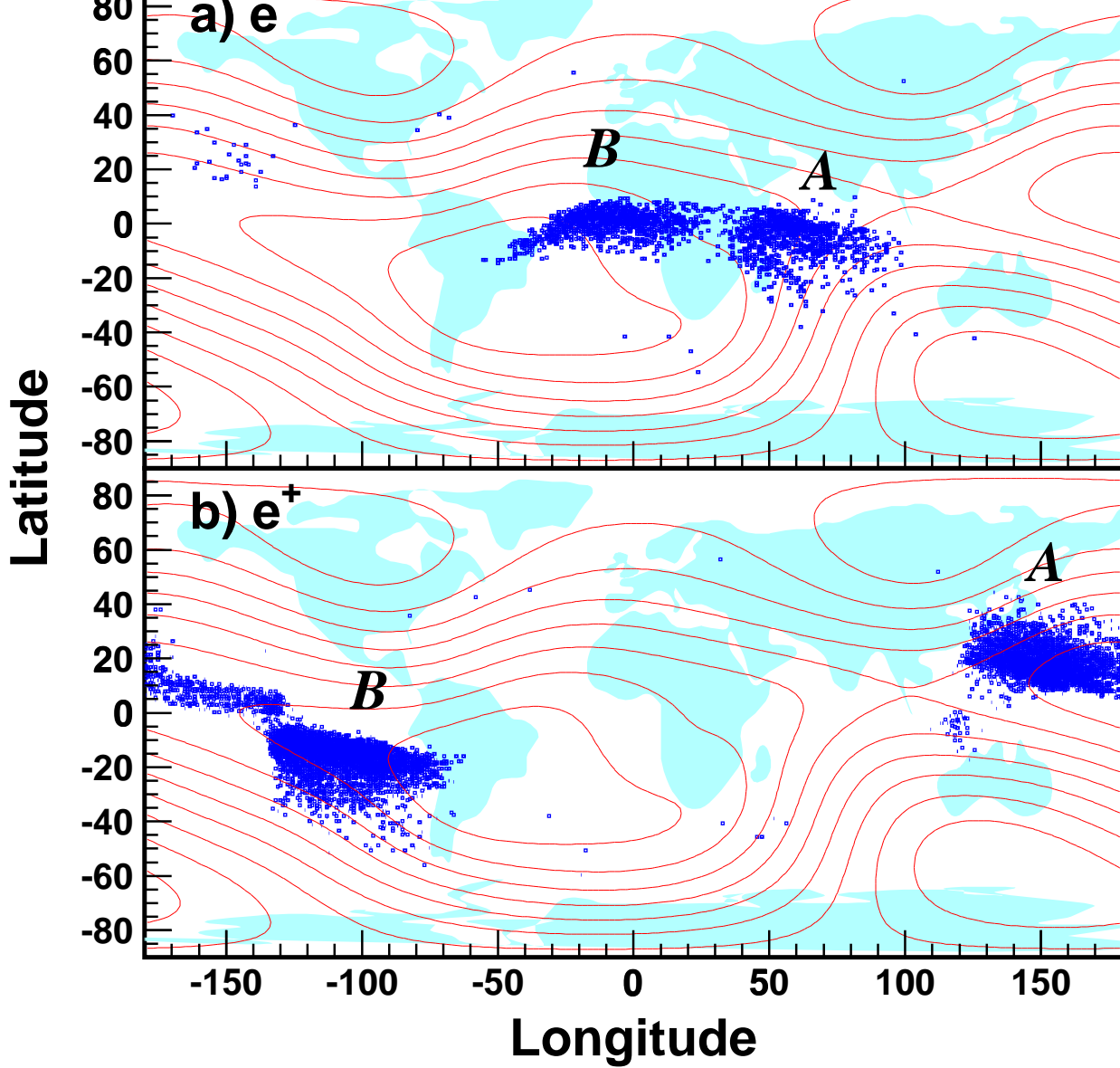


Figure 8: Properties of long-lived second spectra leptons (< 3 GeV): (a) The geographical origin of electrons and (b) positrons. The lines indicate the geomagnetic field contours at 380 km. The regions A and B correspond to the bands A and B marked in Fig. 6. (c) The e^- (full circles) and e^+ (open circles) fluxes integrated over the range 0.2-2.5 GeV as a function of magnetic latitude for zenith, (d) 20° and (e) 45° shuttle attitude.

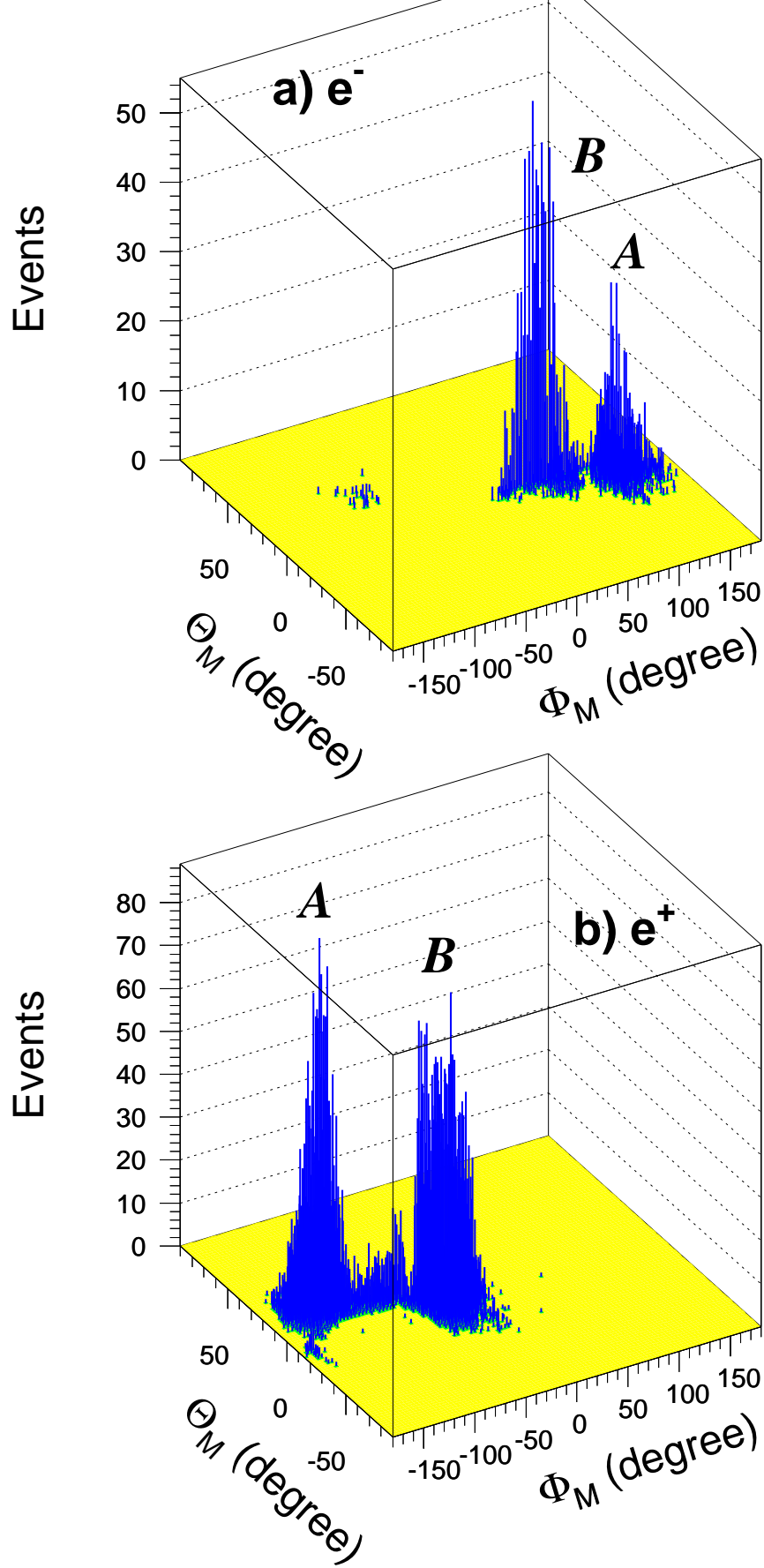


Figure 9: Property of second spectra: The point of origin of long-lived leptons with energies < 3 GeV and $\Theta_M < 0.7$ in geomagnetic coordinates. The regions **A** and **B** correspond to those in Fig. 8 and the bands marked **A** and **B** in Fig. 6.

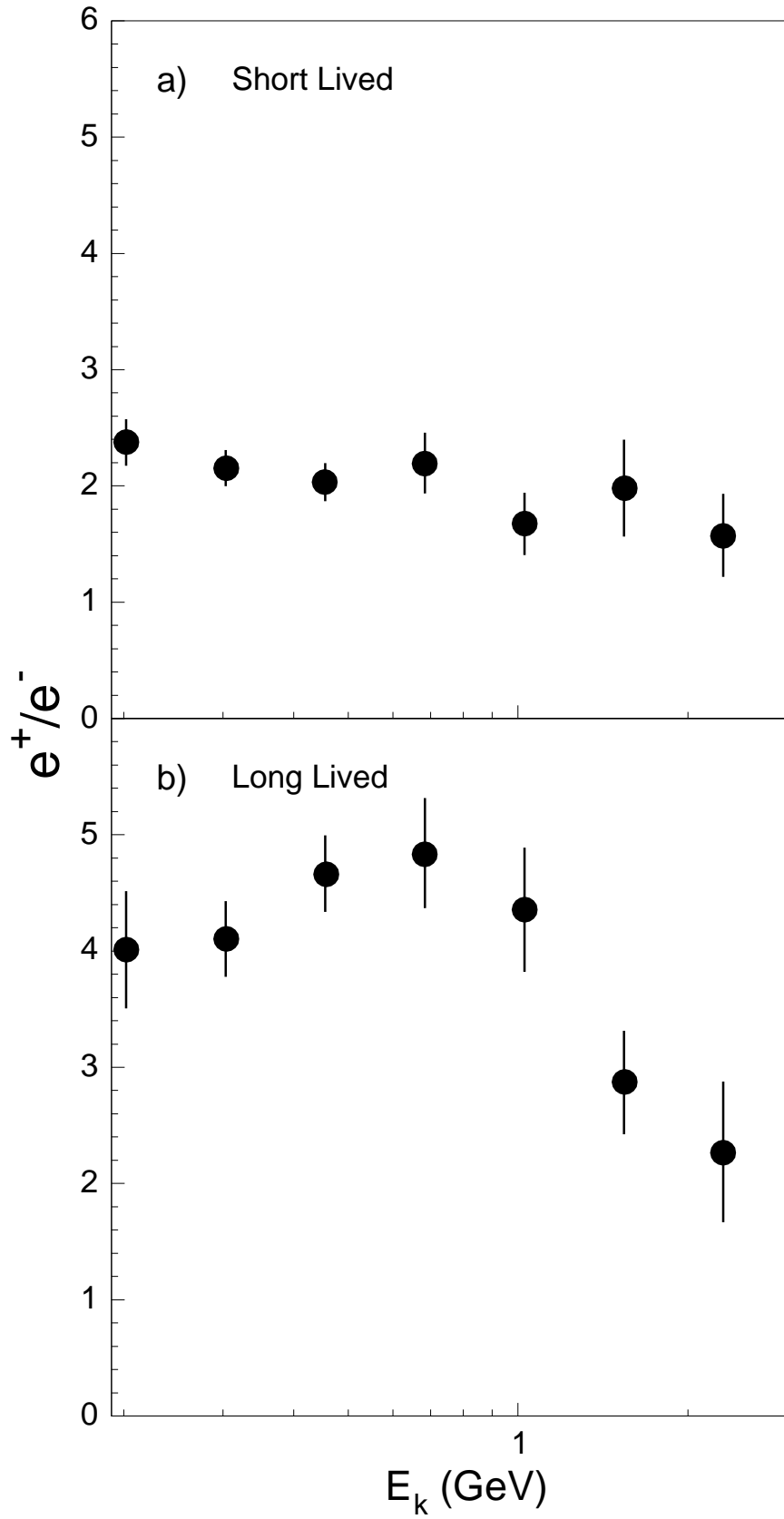


Figure 10: Property of second spectra: The e^+/e^- ratio as a function of energy for (a) short-lived and (b) long-lived particles. Shuttle attitude 0° and $\Theta_M < 0.3$.

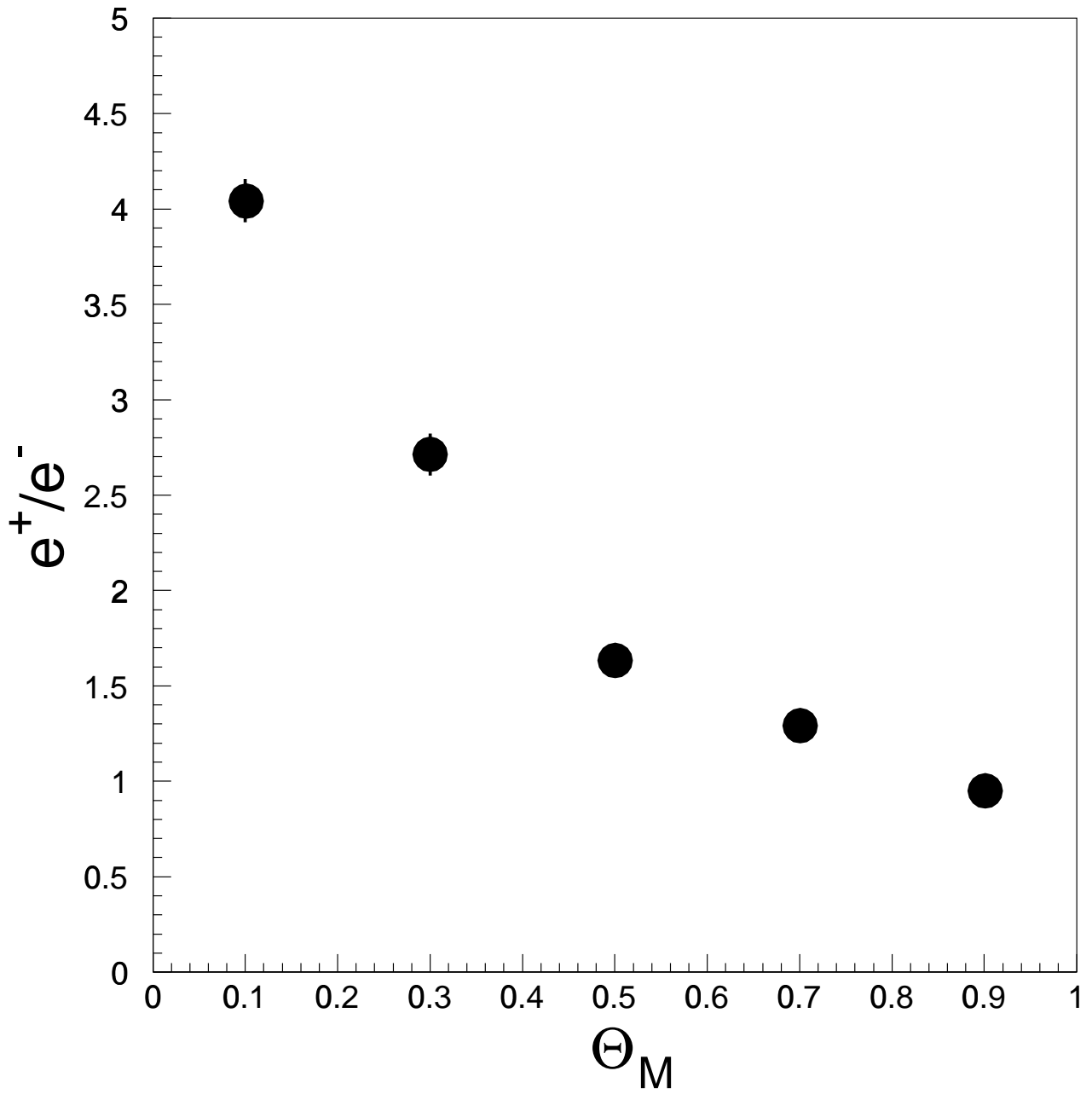


Figure 11: Property of second spectra: The e^+/e^- ratio as a function of magnetic latitude integrated over the range 0.2-2.5 GeV and combined for short-lived and long-lived leptons independent of shuttle attitude.